
This is an electronic reprint of the original article.

This reprint may differ from the original in pagination and typographic detail.

Mack, Iris; Rosta, Kawa; Quliyeva, Ulviyya; Ott, Jennifer; Pasanen, Toni; Vähänissi, Ville; Rad, Zahra Jahanshah; Lehtiö, Juha Pekka; Laukkanen, Pekka; Soldano, Caterina; Savin, Hele

Quantifying the Impact of Al Deposition Method on Underlying Al₂O₃/Si Interface Quality

Published in:

Physica Status Solidi. A: Applications and Materials Science

DOI:

[10.1002/pssa.202200653](https://doi.org/10.1002/pssa.202200653)

Published: 01/10/2023

Document Version

Publisher's PDF, also known as Version of record

Published under the following license:

CC BY

Please cite the original version:

Mack, I., Rosta, K., Quliyeva, U., Ott, J., Pasanen, T., Vähänissi, V., Rad, Z. J., Lehtiö, J. P., Laukkanen, P., Soldano, C., & Savin, H. (2023). Quantifying the Impact of Al Deposition Method on Underlying Al₂O₃/Si Interface Quality. *Physica Status Solidi. A: Applications and Materials Science*, 220(20), Article 2200653. <https://doi.org/10.1002/pssa.202200653>

This material is protected by copyright and other intellectual property rights, and duplication or sale of all or part of any of the repository collections is not permitted, except that material may be duplicated by you for your research use or educational purposes in electronic or print form. You must obtain permission for any other use. Electronic or print copies may not be offered, whether for sale or otherwise to anyone who is not an authorised user.

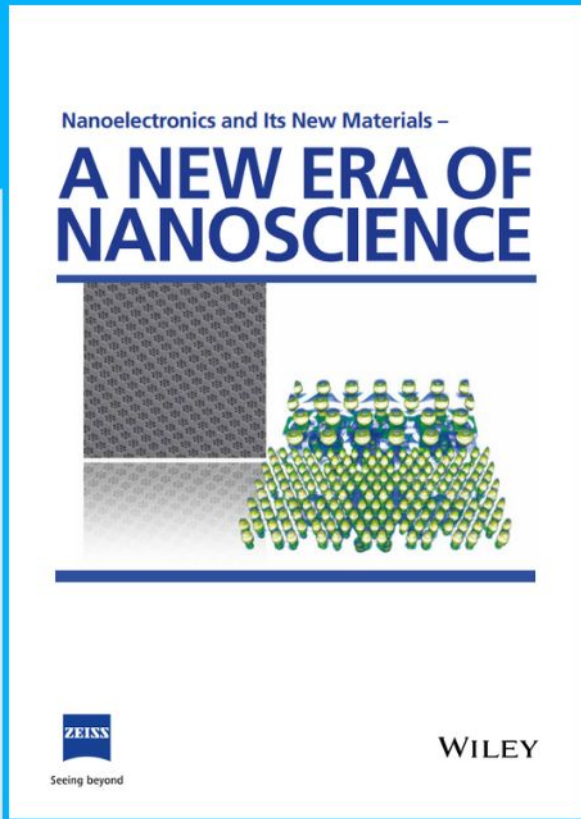


Nanoelectronics and Its New Materials – A NEW ERA OF NANOSCIENCE

Discover the recent advances in electronics research and fundamental nanoscience.

Nanotechnology has become the driving force behind breakthroughs in engineering, materials science, physics, chemistry, and biological sciences. In this compendium, we delve into a wide range of novel applications that highlight recent advances in electronics research and fundamental nanoscience. From surface analysis and defect detection to tailored optical functionality and transparent nanowire electrodes, this eBook covers key topics that will revolutionize the future of electronics.

To get your hands on this valuable resource and unleash the power of nanotechnology, simply download the eBook now. Stay ahead of the curve and embrace the future of electronics with nanoscience as your guide.



Seeing beyond

WILEY

Quantifying the Impact of Al Deposition Method on Underlying $\text{Al}_2\text{O}_3/\text{Si}$ Interface Quality

Iris Mack,* Kawa Rosta, Ulviyya Quliyeva, Jennifer Ott, Toni P. Pasanen, Ville Vähänissi, Zahra Sadat Jahanshah Rad, Juha-Pekka Lehtiö, Pekka Laukkanen, Caterina Soldano, and Hele Savin

Oxide–semiconductor interface quality has often a direct impact on the electrical properties of devices and on their performance. Traditionally, the properties are characterized through metal–oxide–semiconductor (MOS) structures by depositing a metal layer and measuring the capacitance–voltage (C – V) characteristics. However, metal deposition process itself may have an impact on the oxide and the oxide–semiconductor interface. The impact of magnetron sputtering, e-beam evaporation, and thermal evaporation on an $\text{Al}_2\text{O}_3/\text{Si}$ interface is studied, where atomic layer deposited (ALD) Al_2O_3 is used, by MOS C – V and corona oxide characterization of semiconductors (COCOS) measurements. The latter allows characterization of the interface also in its original state before metallization. The results show that sputtering induces significant damage at the underlying $\text{Al}_2\text{O}_3/\text{Si}$ interface as the measured interface defect density D_{it} increases from 10^{11} to $10^{13} \text{ cm}^{-2} \text{ eV}$. Interestingly, sputtering also generates a high density of positive charges Q_{tot} at the interface as the charge changes from -2×10^{12} to $+7 \times 10^{12} \text{ cm}^{-2}$. Thermal evaporation is found to be a softer method, with modest impact on D_{it} and Q_{tot} . Finally, Annealing heals the damage but has also a significant impact on the charge of the film recovering the characteristic negative charge of Al_2O_3 ($\sim -4 \times 10^{12} \text{ cm}^{-2}$).

1. Introduction

Oxide thin films, such as silicon oxide (SiO_2), aluminum oxide (Al_2O_3), or hafnium oxide (HfO_2), are of great importance for the integrated circuit industry, as well as in many optoelectronic applications due to their intrinsic properties. For instance, they are often used as optical or insulation layers in devices^[1–4] and they are capable of providing a high-quality interface with neighboring materials such as silicon. In electrical devices, a high-quality interface possesses a low recombination rate, in other words, oxide films can reduce the recombination probability of charge carriers at the silicon surface, which is often referred to as surface passivation.^[5–12]

Efficient surface passivation by oxide films is based on two different mechanisms. The first mechanism is to decrease the interface defect density (D_{it}) at the oxide–semiconductor interface, by reducing the amount of dangling bonds.^[13] For


this, SiO_2 is one of the most widely used oxide films in Si-based devices.^[14] The second one is the so-called field-effect mechanism, which significantly reduces the electron or hole concentration at the oxide–semiconductor interface by means of a built-in electric field.^[15] This internal electric field can be obtained by the presence of an increased electrical charge density (Q_{tot}) at the oxide–semiconductor interface. A well-known material for this is Al_2O_3 , which has a high negative fixed charge density. Combining a high Q_{tot} with a low D_{it} would be ideal, which can be obtained with $\text{Al}_2\text{O}_3/\text{SiO}_x/\text{Si}$ interface.^[7] In this context, it is important to point out that Al_2O_3 deposited by atomic layer deposition (ALD), possesses negative fixed charges regardless of whether the layer is deposited on H-terminated silicon or on silicon covered with a very thin SiO_2 layer.^[15–17] Indeed, several studies have shown that an interfacial SiO_2 layer is formed between the Al_2O_3 film and silicon, where the negative fixed charges are located.^[7–23]

In addition to dielectric layers, semiconductor devices typically contain also metal layers, for example, to realize electric contacts. One main challenge is to maintain the high quality of the oxide and the oxide–semiconductor interface, even after metal layer deposition. Different metallization methods, for example,

I. Mack, K. Rosta, U. Quliyeva, J. Ott, T. P. Pasanen, V. Vähänissi, C. Soldano, H. Savin
Department of Electronics and Nanoengineering
Aalto University
Tietotie 3, FI-02150 Espoo, Finland
E-mail: macki@mailbox.org

J. Ott
Helsinki Institute of Physics
University of Helsinki
Gustaf Hållströminkatu 2, FI-00014 Helsinki, Finland

Z. S. Jahanshah Rad, J.-P. Lehtiö, P. Laukkanen
Department of Physics and Astronomy
University of Turku
Vesilinnantie 5, FI-20014 Turku, Finland

 The ORCID identification number(s) for the author(s) of this article can be found under <https://doi.org/10.1002/pssa.202200653>.

© 2023 The Authors. physica status solidi (a) applications and materials science published by Wiley-VCH GmbH. This is an open access article under the terms of the Creative Commons Attribution License, which permits use, distribution and reproduction in any medium, provided the original work is properly cited.

DOI: 10.1002/pssa.202200653

physical vapor deposition (PVD) techniques, such as magnetron sputtering, electron beam (*e*-beam) evaporation, and thermal evaporation, introduce different types and quantity of damage. For instance, during magnetron sputtering, the substrate is bombarded by a variety of particles, such as low-energy argon ions from the plasma, sputtered metal atoms, or high-energetic electrons that generate electrically active damage in the underlying layers.^[24–27] In *e*-beam evaporation, low-energy X-rays from the electrons evaporating the source material introduce defects in the metal–dielectric–semiconductor stack.^[28–30] In comparison, thermal evaporation is considered as the softest PVD method for depositing metal contacts on oxide–semiconductor structures. Based on the above, it is generally known that the damage introduced by the described metallization methods is different; however, investigations on their effect on D_{it} and Q_{tot} are rather limited, as deep-level transient spectroscopy (DLTS) and electron-beam-induced current (EBIC) techniques are used to determine the nature of the defects and their influence on recombination.^[24,26–31]

Despite the possible damage to the underlying layers, metal layers can also be used to improve interface quality. A known and widely used method to reduce the defect density at the oxide–silicon interface, when using aluminum (Al) as metal contact, is to perform a postdeposition anneal after the Al deposition.^[5,32–34] This method was named Alneal by B. Deal as a combination of Aluminum and anneal, when first investigated in 1969.^[34] The effect of Alnealing was further investigated, for example, by Kerr and Cuevas through photoconductance^[14] and by W. Eades through DLTS.^[35] Furthermore, Collett et al. studied that the effect can be enhanced by adding different ions at the Si–SiO₂ interface.^[32] Setälä et al. observed that Alnealing can have an influence on the passivation performance of Al₂O₃.^[36] Although it is known that Alnealing can efficiently reduce dangling bonds at the oxide–Si interface in some configurations, it has remained unresolved, if the Al deposition method influences the efficiency of Alnealing.

The electrical quality of the oxide–semiconductor interface is typically evaluated by capacitance–voltage (*C*–*V*) measurements, in which the semiconductor/oxide structure is integrated in a metal–oxide–semiconductor (MOS) capacitor by adding a patterned metal layer on top.^[37] However, this *C*–*V* technique has also some drawbacks: it is necessary to prepare separate samples that require metal contacts, the contact area needs to be accurately geometrically defined, as this influences the fitting of the data, and to get quantitative results either both true low- and high-frequency data is needed or when using Terman method true high-frequency and theoretical *C*–*V* curves are needed.^[38] Moreover, as the metal deposition is suspected to have an influence on the properties of the oxide–semiconductor interface, as stated earlier, the fabrication of MOS capacitors for the *C*–*V* measurements may influence the results and the measurement data does not necessarily represent the original interface. These limitations can be overcome using the corona oxide characterization of semiconductors (COCOS) method that is based on measuring the surface contact potential via Kelvin probe.^[39,40] COCOS provides a great advantage, since it is a noninvasive contactless technique, it allows to characterize the interface without the need of a metal electrode, and it is independent of the sample

geometry. Thus, COCOS represents an ideal tool to determine the nature of the oxide–semiconductor interface prior to metal deposition and throughout different device fabrication steps.^[39,41]

In this article, we investigate the Al₂O₃/Si interface properties and how they are affected by different metallization methods. We present a quantitative comparison of the influence of magnetron sputtering, *e*-beam evaporation, and thermal evaporation. Aluminum is chosen as metal due to its common use in silicon-based device fabrication. We investigate Q_{tot} and D_{it} by transforming the Al₂O₃/Si stack into MOS capacitors for *C*–*V* measurements and by measuring it directly with the COCOS method. By comparing the results obtained by these two methods, we determine the influence the aluminum deposition method has on the electrical properties of the interface. Furthermore, we assess the advantages and disadvantages of the two measuring methods for the characterization of the oxide–Si interface quality. In addition, we investigate with COCOS what impact the Alneal process in comparison to annealing without Al on top has on D_{it} and Q_{tot} and whether the used metallization method has any impact on the efficiency of Alnealing.

2. Experimental Section

The experiments were performed on single-side polished p-type float-zone silicon wafers with $3.00 \pm 2.00 \Omega \text{ cm}$ resistivity. The wafer thickness was $280 \pm 15 \mu\text{m}$. The main process (blue) and characterization steps (green and yellow) are schematically presented in Figure 1.

The wafers were first cleaned using the SC1 and SC2 cleaning (SC1 standard clean: NH₄OH:H₂O₂:H₂O (1:1:5) at 80 °C for 10 min. SC2 standard clean: HCl:H₂O₂:H₂O (1:1:5) at 80 °C for 15 min) with an intermediate HF dip (HF:DIW 1:50). Then, ~22 nm-thick Al₂O₃ films were deposited by ALD in a Beneq TFS-500 system, at a temperature of 200 °C, using trimethylaluminum (TMA) and water (H₂O) as precursors. Subsequently, the samples were annealed in a Centrotherm

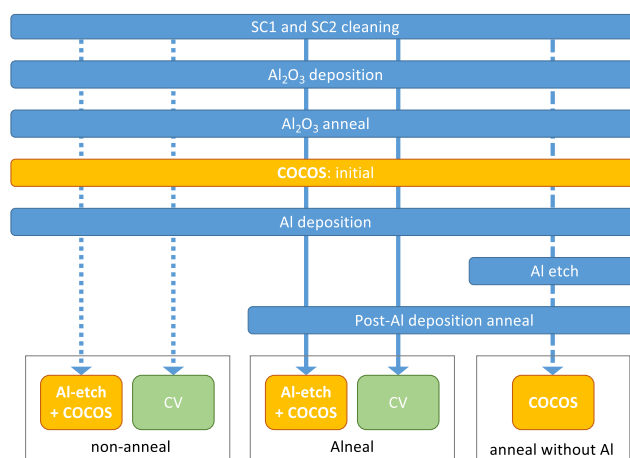


Figure 1. Schematic of the processing steps, as well as the corresponding characterization steps. The yellow and green boxes indicate the COCOS and conventional *C*–*V* samples compared in this study.

furnace at 400 °C for 30 min in N₂ atmosphere (labeled “Al₂O₃ anneal” in Figure 1) to activate the surface passivation.

Afterward, an aluminum layer was deposited on the Al₂O₃ film at room temperature by three different deposition methods: magnetron sputtering, *e*-beam evaporation, or thermal evaporation. The sputtered films were deposited at a base pressure of 2.3×10^{-7} mbar from a pure Al target driven at 2 kW DC power and with an argon (Ar) gas flow of 22 sccm. For *e*-beam evaporation, the base pressure was 5×10^{-7} mbar, the acceleration voltage 10 kV, and the beam current 20 mA. Thermal evaporation was performed at a base pressure of 1.2×10^{-6} mbar. The overall thickness of the Al films was around 200 nm for all deposition methods.

Following the Al deposition, most of the wafers were annealed at 400 °C for 30 min in N₂ atmosphere (“post-Al deposition anneal” in Figure 1). In other words, these wafers experienced the annealing step, which is commonly known as Alneal.^[34] As a reference, in some of the samples, the Al layer was removed prior to exposing them to the same post-Al deposition anneal (dashed path in Figure 1). In these wafers, Al was removed using phosphoric acid-based Al etch mixture and the etching time was chosen in such a way that only aluminum and not the underlying Al₂O₃ was removed.

Two methods were used to characterize the interface properties: conventional contact *C*–*V* measurements for MOS capacitor samples (green squares in Figure 1) and COCOS measurements for parallel produced samples, where the Al was removed before measuring (yellow squares). For the MOS capacitors, Al was deposited through a shadow mask with varying contact sizes from 0.008 to 7.1 mm². The backside contact was realized by applying silver paste on the rear side of the wafer. *C*–*V* curves of the MOS capacitors were measured with a probe station^[37] at 100 kHz and the obtained results for a single contact size were compared.

The COCOS method was based on external corona charge deposition on top of the oxide layer.^[39,40] Therefore, it was necessary to remove the Al films prior to the COCOS measurements using the same Al etch as described earlier. The corona charge

was applied in small increments, from inversion to accumulation, and after each step the resulting contact potential difference was measured both in complete darkness and under strong illumination.^[39] From these measurements, electrical parameters of the oxide–semiconductor stack such as flatband voltage (*V*_{FB}), *Q*_{tot}, and *D*_{it} were calculated. Every sample was measured in three different locations, and the range of variation is presented in each figure as error bars.

Conventional *C*–*V* measurements and COCOS were carried out on two sample types for comparison purposes: right after Al deposition (dotted path in Figure 1) and after the following post-Al anneal (solid path). As mentioned above, for the COCOS samples, the Al layer was removed prior to the measurements. The oxide charge (*Q*_{ox}) in MOS capacitors, which is assumed to be located at the oxide–semiconductor interface, was determined by measuring the *C*–*V* curves from inversion to accumulation for p-type substrates. COCOS measurements were also performed from inversion to accumulation. The determined *Q*_{tot} was the total charge needed to achieve flatband condition starting from the initial state of the oxide–semiconductor system.

Additionally, COCOS allows characterization of the oxide–silicon interface right after Al₂O₃ deposition (i.e., prior to the Al deposition) as well as of another sample type, where the Al layer was etched away before the sample obtained the post-Al deposition anneal (dashed line). These measurements were not possible with conventional MOS capacitors due to the absence of the metal layer.

3. Results and Discussion

3.1. Impact of Metal Deposition Method

The impact of the different metallization methods on the Al₂O₃/Si interface was investigated by measuring the *C*–*V* characteristics from MOS capacitor samples, as well as by quantifying *Q*_{tot} and *D*_{it} from the corresponding COCOS samples. Figure 2a presents the *C*–*V* curves for MOS capacitors metallized by

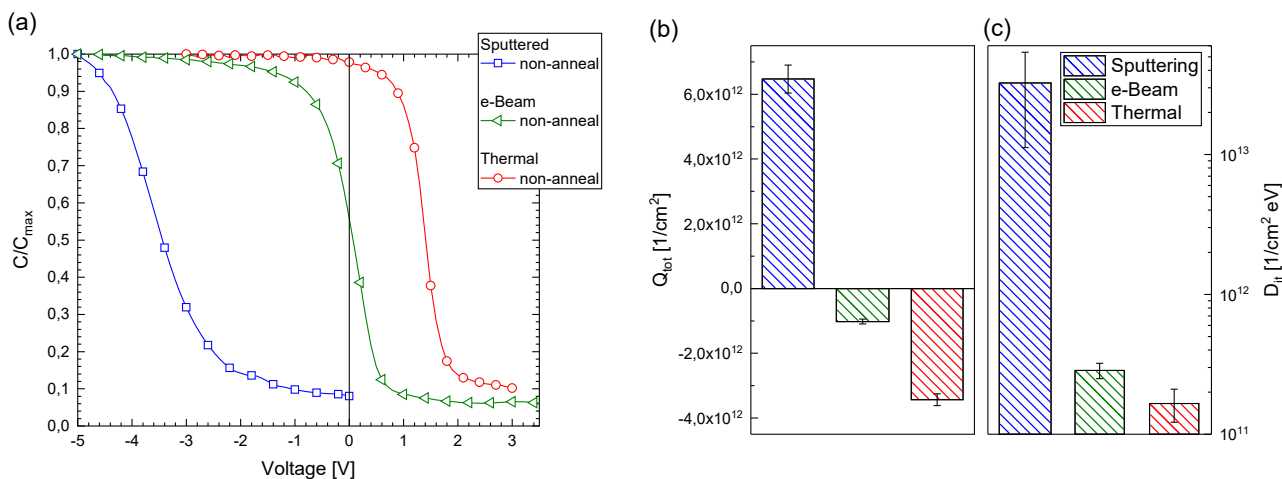


Figure 2. Influence of different metallization methods for Al without postmetallization anneal on a) the *C*–*V* characteristics (*CV* non-anneal) of MOS capacitors, and b) the total charge *Q*_{tot} and c) the interface defect density *D*_{it} at midgap by COCOS measurements (COCOS non-anneal).

sputter deposition, e -beam, and thermal evaporation directly after the Al deposition (CV non-anneal). The curves were normalized to the largest value of each curve at accumulation.

The fast decrease of the capacitance in the C - V characteristics corresponds to the so-called flatband voltage (V_{FB}), at which the energy bands of the substrate are becoming flat at the Al_2O_3/Si interface. The difference of V_{FB} shows that the deposition method has a large impact on the amount and type of charge at the interface. The negative flatband voltage that is observed for the sputtered sample (blue, open squares) indicates that after Al sputtering, there are positive charges at the oxide-semiconductor interface. After metallization by e -beam evaporation (green, open triangles), V_{FB} reaches near-zero voltage. This reveals that there is only little-to-no charge at the interface. However, after thermal evaporation (red, open circles), negative charge is found at the interface, which is expected due to the negatively charged Al_2O_3 . From the flatband voltage shift, the oxide charge Q_{ox} at the oxide-semiconductor interface can be determined through

$$Q_{ox} = -C_{ox} \Delta V_{FB} \quad (1)$$

where C_{ox} is the capacitance of the oxide in accumulation and ΔV_{FB} the flatband voltage shift from neutrality (0 V) due to total oxide and interface charges.^[42] For the investigated contacts, the derived Q_{ox} is for the sputtered contacts (blue) $(+10.9 \pm 5.9) \times 10^{12} \text{ cm}^{-2}$, whereas $(-0.2 \pm 0.3) \times 10^{12} \text{ cm}^{-2}$ for the e -beam evaporated (green) and $(-2.1 \pm 2.0) \times 10^{12} \text{ cm}^{-2}$ for thermal evaporated (red).

In Figure 2a it is also visible that the slope of the normalized C - V curves (i.e., the interface defect density) for the three metallization methods is different. The curves indicate that sputtering generates the highest D_{it} (lowest slope), while thermal evaporation seems to be the least damaging leading to the lowest D_{it} (steepest slope).

COCOS measurements were performed on sister samples to determine Q_{tot} and D_{it} . Q_{tot} is defined as

$$Q_{tot} = Q_{it}^{FB} + Q_{surf} + Q_{ox} \quad (2)$$

where Q_{it}^{FB} is the interface-trapped charge at flatband condition, Q_{surf} is any initial charge on the surface of the oxide, and Q_{ox} is the oxide charge. The results for Q_{tot} and D_{it} are given in Figure 2b,c, respectively. For the sputtered contact (blue), the total charge Q_{tot} is $(+6.5 \pm 0.5) \times 10^{12} \text{ cm}^{-2}$, whereas it is $(-1.0 \pm 0.1) \times 10^{12} \text{ cm}^{-2}$ for e -beam (green) and $(-3.4 \pm 0.2) \times 10^{12} \text{ cm}^{-2}$ for thermally (red) evaporated aluminum. These values show the same trends as the above-described qualitative MOS and Q_{ox} results. They confirm that there are a large number of positive charges at the oxide-semiconductor interface for the sputtered and of negative charges for the thermal evaporated samples. While D_{it} was more difficult to interpret from the C - V MOS results, Figure 2c indicates significant difference in the D_{it} at the Al_2O_3/Si interface for the three metallization methods. For the sputtered contact (blue), the defect density $((3.3 \pm 2.2) \times 10^{13} \text{ cm}^{-2} \text{ eV})$ is more than two orders of magnitude higher than for the evaporated contacts: e -beam (green) $(2.9 \pm 0.4) \times 10^{11} \text{ cm}^{-2} \text{ eV}$ and thermal (red) $(1.7 \pm 0.5) \times 10^{11} \text{ cm}^{-2} \text{ eV}$.

Taking into account the definition of Q_{tot} given in Equation (2), the differences in the determined values for Q_{tot} and Q_{ox} can be explained. The values for Q_{tot} are always larger, as they also include the interface-trapped charge (Q_{it}^{FB}) and the initial surface charges (Q_{surf}). Q_{surf} should be equal for all samples, as the preparation was identical. On the other hand, Q_{it}^{FB} should strongly depend on the deposition method, as the above determined D_{it} strongly varies for the different deposition methods. This is in line with the obtained values for Q_{tot} and Q_{ox} , as the difference between the two and the corresponding D_{it} is larger for the sputtered, than for the evaporated samples.

The obtained differences for the different metal deposition methods were expected to some extent, as it is known that, for example, during sputtering, the sample is bombarded with a variety of particles that can cause damage,^[24–26] but the quantified amount is new. As reported by B. Deal, ionizing radiation of the plasma during sputter deposition generates electron-hole pairs in the oxide, which result in induced positive charge due to trapped holes.^[34] However, it is still remarkable that sputtering generates D_{it} above $1 \times 10^{13} \text{ cm}^{-2} \text{ eV}$, as such high defect densities are not tolerable in devices.^[43] This explains, why, curing steps such as annealing are in general introduced in fabrication processes.

3.2. Impact of Alneal

In SiO_2/Si -based device fabrication, it is common to intentionally cover the oxide-semiconductor surfaces with aluminum and perform a post-Al anneal (Alneal) to improve the quality of the SiO_2/Si interface.^[14,32,34] To investigate the effect of the Alneal on the Al_2O_3/Si interface, we compared in Figure 3a the C - V results of Alnealed samples (CV Alneal: filled symbols) to the above-described as-deposited samples (CV non-anneal: open symbols).

Upon Alneal, all C - V curves shifted toward similar flatband voltages; therefore, the major differences for the different deposition methods discussed earlier have been reduced. In particular, we observed that 1) positive charges induced by Al sputter deposition are cancelled out and actually become negative and 2) the (negative) charge density of the e -beam evaporated sample is further increased. In addition, differences in the slopes have been reduced, indicating that the Alneal process is effective in recovering most of the damage (possibly all) induced by metal deposition and in reducing any intrinsic difference derived by the metal deposition method. These changes in the C - V curves also reflect in the values of Q_{ox} determined from them, leading to same values within the measurement accuracy of the method. In numerical numbers, the total charge after Alnealing for Al deposited by sputtering, e -beam, or thermal evaporation is $(-4.0 \pm 3.1) \times 10^{12}$, $(-1.8 \pm 2.0) \times 10^{12}$, and $(-3.1 \pm 2.6) \times 10^{12} \text{ cm}^{-2}$, respectively.

The changes observed with the C - V curves are also reflected in the values of Q_{tot} and D_{it} measured with COCOS. By comparing the filled (COCOS Alneal) with the ruled (COCOS non-anneal) bars in Figure 3b, one can see that Q_{tot} at the oxide-silicon interface after Alnealing is about

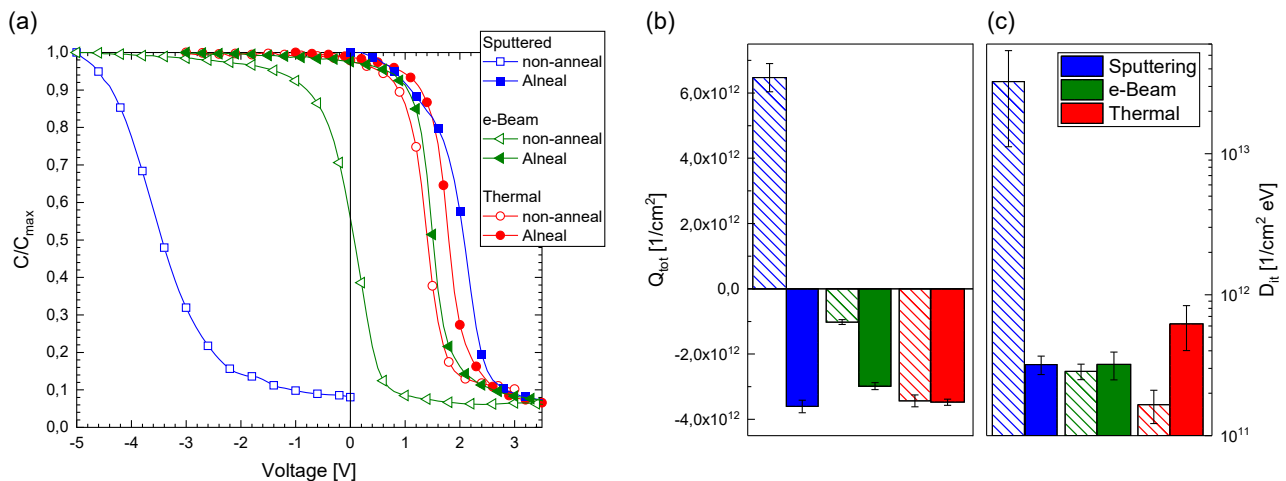


Figure 3. Influence of post-Al deposition anneal (Alneal) on the oxide-semiconductor structure on a) the C–V characteristics for different metallization methods before (open symbols, CV non-anneal) and after Alnealing (filled symbols, CV Alneal). b) Total charge Q_{tot} and c) interface defect density D_{it} at midgap with respect to the influence of metallization method for non-anneal (ruled, COCOS non-anneal) and Alneal (filled, COCOS Alneal) samples.

$(-3.3 \pm 0.3) \times 10^{12} \text{ cm}^{-2}$ for all samples, independently of the prior metallization method. Upon Alneal, also the D_{it} values are all within the same order of magnitude, independently of the deposition method, as shown in Figure 3c. This is in agreement with the changes in the slopes observed in the C–V curves in Figure 3a and the derived values given earlier. We note a non-negligible increase of the interface defect density after Alneal treatment for the thermally evaporated samples measured by COCOS; this unexpected change is currently under investigation.

Our experimental results demonstrate that even a large damage induced by metallization (as in the case of sputtering) can be cured by postmetallization annealing (Alneal) of the samples. According to B. Deal, the reduction of the positive charge due to trapped holes, induced by the plasma during sputtering, can be annealed out at 350 °C or higher.^[34] The curing effect for D_{it} by Alnealing was also expected to a certain extent, as reductions for Al/SiO₂/Si-based structures are reported in earlier works.^[5,15,19,34,43] Furthermore, a similar reduction in D_{it} of Al/Al₂O₃/Si MOS structures with sputtered Al films was recently reported by Khosla et al.^[44] The improvement could be explained by hydrogen species, which diffuse to the oxide-semiconductor interface and terminate the dangling bonds at the interface, as done for SiO₂/Si interfaces.^[5,32,34,43,45] Those hydrogen species form, when water, introduced to the system as traces of moisture, reacts with the Al during annealing. Water traces are introduced at different points throughout the process and can therefore be at the Al-oxide interface, in the Al or on the Al surface.^[5,32,34,43,45] The beneficial impact of Alnealing has also been observed, when using SiO₂/TiO₂ as oxide in metal-oxide-semiconductor thin-film stacks,^[32,33] which supports the assumption that the same theory could also be applied for our Al₂O₃/Si interface. In summary, the above-discussed experiments have shown that Alnealing not only reduces the D_{it} at the oxide-semiconductor interface, but also significantly changes the charge density Q_{tot} .

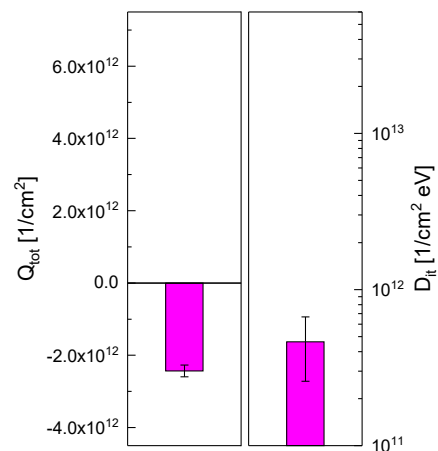


Figure 4. Initial values for total charge Q_{tot} and interface defect density D_{it} at midgap for 22 nm-thick Al₂O₃ measured directly after oxide deposition and Al₂O₃ anneal (COCOS initial).

3.3. Initial Properties of the Al₂O₃/Si Interface

We further used COCOS characterization to study the initial stage (before metallization) of the studied Al₂O₃/Si interface right after activating the surface passivation (Al₂O₃ anneal). The measured values for Q_{tot} $((-2.4 \pm 0.2) \times 10^{12} \text{ cm}^{-2})$ and D_{it} $((4.6 \pm 2.1) \times 10^{11} \text{ cm}^{-2} \text{ eV})$ (COCOS initial) are presented in Figure 4, and are consistent with literature.^[15,20–22,46]

When comparing with the corresponding Q_{tot} after metal deposition (Figure 2), it becomes clear that the initial value is intermediate between the values after *e*-beam and thermal evaporation. On the other hand, if we compare the initial D_{it} to D_{it} of thermal and *e*-beam evaporated samples, it seems that deposition of Al by evaporation actually improves the Al₂O₃/Si interface. When comparing the obtained initial values in Figure 4 to the values given in Figure 3 for the Alnealed samples, it becomes

visible that by Alnealing of the oxide–semiconductor interface, the characteristic negative charge of Al_2O_3 can be increased by $\Delta Q_{\text{tot}} \approx 1 \times 10^{12} \text{ cm}^{-2}$ and that D_{it} changes back to the value of the initial stage. A similar increase of Q_{tot} was also observed by Setälä et al. for n-type Si-wafers, and they also observe a decrease for D_{it} after Alnealing.^[36]

3.4. Impact of Presence of Al Layer During Alneal

In the sections above, we have demonstrated that Alnealing is effective for healing damage induced by metal deposition. However, it remains unclear if this effect is a direct consequence of the annealing process itself or possibly related to the physical presence of the Al layer during annealing. COCOS allows us to gain a deeper insight in this subject. Samples which have undergone the same anneal process (as Alnealed samples) after the aluminum layer was removed (COCOS anneal) were prepared and characterized. Figure 5 shows the effect of the annealing process with and without Al layer, filled and dotted bars, respectively. Our results suggest that the value of Q_{tot} depends on the presence of the Al film during the thermal treatment, independently of the deposition method, the absolute value of Q_{tot} is consistently smaller after annealing (COCOS anneal – dotted) than after Alnealing (COCOS Alneal – filled). In addition, when comparing Q_{tot} after annealing with the initial values (Figure 4), it is clear that the values after annealing without Al (COCOS anneal) are closer to the initial value than after Alnealing. Therefore, we can conclude that the annealing process by itself (at 400 °C in pure N_2 atmosphere) has a positive effect, as it reduces the positive charges induced by Al deposition. This is consistent with work reported by B. Deal, as he has reported earlier that plasma-induced positive charge can be cured by annealing at temperatures higher than 350 °C.^[34] Furthermore, the investigation of Q_{tot} shows that it is enough to anneal after metal deposition to regain the original strength of field-effect passivation. On the other hand, for further improvement in field-effect

passivation, it is desirable to increase negative charges at the interface; thus, Alnealing is in this case the preferred choice.

Regarding the interface defect density, the effect is quite surprisingly overall less clear (Figure 5b) even though Alnealing was introduced to improve the oxide–semiconductor interface. The results indicate negligible differences between Alnealing and annealing for the sputter-deposited (blue) or *e*-beam-evaporated (green) samples. In the case of thermally evaporated films (red), a more pronounced difference in D_{it} is observed. This behavior is currently investigated further. We also note that independently of the annealing process (i.e., with or without Al on the surface), D_{it} values are smaller than the measured initial value (Figure 4), except for the thermally evaporated Alnealed sample (red filled). This could indicate that in the case of ALD-deposited Al_2O_3 , the dissociation of water by Al during the Alnealing is less important, as the Al_2O_3 layer already contains a remarkable amount of hydrogen that can heal dangling bonds at the interface.^[6]

3.5. Discussion and Summary

The systematic investigation presented here demonstrates that the metallization method influences the characteristics of oxide–semiconductor interfaces. Table 1 summarizes the obtained values for Q_{tot} and D_{it} for the different deposition methods and the various processes measured with the COCOS method.

By preparing identical samples, where only the Al deposition method was varied, we confirmed that sputtering indeed induces a high amount of damage at the interface, which can be seen in the high D_{it} value. Furthermore, thermal evaporation can be considered as the least damaging method among the three PVD deposition techniques investigated. On the other hand, when comparing the non-anneal values for Q_{tot} with the initial values obtained for the pure $\text{Al}_2\text{O}_3/\text{Si}$ interface (before any Al deposition), it becomes clear that the total negative charge at the oxide–semiconductor interface, which is needed for the field-effect passivation, is reduced during the Al deposition by sputtering or *e*-beam evaporation, whereas it is increased during thermal evaporation.

The exact generation of the damage at the interface was not investigated in this article, if it is only the bombardment by different species during metal deposition, as suggested in literature,^[24–26,28–30] or if it is the induced temperature during the metal deposition. Nevertheless, the just described factors should have an influence on the absolute values, especially if

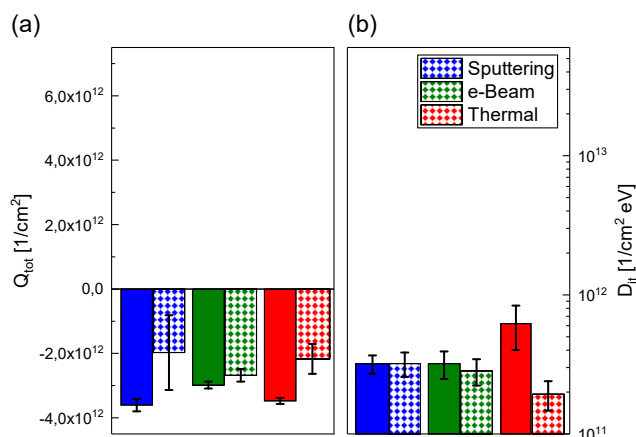


Figure 5. Effect of postmetal deposition anneal on samples with and without the aluminum layer. a) Q_{tot} and b) D_{it} after postmetal annealing with aluminum (filled: COCOS Alneal) or without aluminum (dotted: COCOS anneal) on the oxide measured by the COCOS method.

Table 1. Summary of total charge Q_{tot} and interface defect density D_{it} for Al_2O_3 -passivated samples after Al deposition via different methods and anneal processes, as investigated by COCOS method. (* This unexpected value is currently under investigation).

	Q_{tot} [10^{12} cm^{-2}]			D_{it} [$10^{11} \text{ eV}^{-1} \text{ cm}^{-2}$]		
	Sputtering	<i>e</i> -beam	Thermal	Sputtering	<i>e</i> -beam	Thermal
Initial (before Al)		–2.4			4.6	
Non-anneal	6.5	–1.0	–3.4	325.0	2.9	1.7
Alneal	–3.6	–3.0	–3.5	3.2	3.2	6.2*
Anneal without Al	–2.0	–2.7	–2.2	3.2	2.8	1.9

different metal contacts are used, as they should depend on the deposition parameters needed, such as power or pressure. This is further supported by the findings of Cho et al., who investigated the influence of process parameters when depositing a thin film by PECVD on top of Al_2O_3 .^[47] Furthermore, we have seen that for thin (12 nm) Al_2O_3 the changes to Q_{tot} are more pronounced, although the tendencies are similar to those of the discussed thicker (22 nm) Al_2O_3 films (results not presented in this article). Therefore, we think that we can assume that the observed phenomena could also be valid for other oxide–semiconductor interfaces, independent of the used materials.

Furthermore, the changes in Q_{tot} and D_{it} in comparison to their initial values can be influenced by annealing without metal or by Alnealing the oxide–Si interface. When only D_{it} is considered, there is no clear preference for annealing or Alnealing, as the differences are in the range of measurement fluctuations. On the contrary, if a high field effect, that is, a high Q_{tot} , is needed, which is even higher than the initial values, it is necessary to perform an Alneal.

Finally, the investigated influence of different metal deposition methods and the different annealing steps on the oxide–semiconductor interface show that the COCOS characterization possesses several benefits in comparison to using the second harmonic generation (SHG) method^[48] or the traditional investigation by measuring C – V curves of MOS capacitors. First, there is no requirement to precisely know the sample geometry (SHG and MOS). Second, there is no need for electrical contacts (MOS). Instead, COCOS opens the possibility to assess the initial stage and therefore the influence of the contact deposition, as well as the impact of metal-free annealing in comparison to Alnealing of the oxide–semiconductor interface. Furthermore, reducing the number of process steps required has the advantage of limiting possible contamination.

4. Conclusion

In this article, we systematically investigated the influence of different metallization methods on the characteristics of oxide–semiconductor interfaces by two different methods: on the device level by measuring C – V curves of MOS capacitors and on film level through the COCOS method. By measuring Q_{tot} and D_{it} , we have shown that magnetron sputtering introduces significant damage, whereas thermal evaporation is the softest method. Furthermore, we have demonstrated that Alnealing heals the damage as expected but rather surprisingly it has also a significant impact on the charge of the film recovering the characteristic high negative charge of Al_2O_3 . However, if the anneal is performed without the presence of aluminum, Q_{tot} is not fully recovered whereas D_{it} returns to the original level. In summary, when characterizing interfaces and/or designing devices, one should keep in mind that metallization methods as well as annealing processes can affect the interface quality.

Acknowledgements

I.M. and K.R. contributed equally to this study and should be considered joint first authors. The authors would like to thank Business Finland in the framework of the Silicon Surface Processing Commercialization

(SISUpoco) project, the Academy of Finland (#331313, #338974), and Tandem Industry Academia funding from the Finnish Research Impact Foundation for financial support of the research. The work is related to the Flagship on Photonics Research and Innovation “PREIN” funded by Academy of Finland. Furthermore, the authors acknowledge the provision of facilities by Aalto University and VTT Technical Research Centre of Finland at OtaNano - Micronova Nanofabrication Centre and the technical support by Aalto's Nanofab team.

Conflict of Interest

The authors declare no conflict of interest.

Data Availability Statement

The data that support the findings of this study are available from the corresponding author upon reasonable request.

Keywords

corona oxide characterization of semiconductors, defect density, interface quality, metal–oxide–semiconductors, physical vapor deposition, thin-film metallization

Received: September 19, 2022

Revised: July 23, 2023

Published online:

- [1] J. Kim, J. Park, J. H. Hong, S. J. Choi, G. H. Kang, G. J. Yu, N. S. Kim, H.-E. Song, *J. Electroceram.* **2013**, *300*, 41.
- [2] X. Yu, T. J. Marks, A. Facchetti, *Nat. Mater.* **2016**, *150*, 383.
- [3] A. Prasad, S. Balakrishnan, S. K. Jain, G. C. Jain, *J. Electrochem. Soc.* **1982**, *1290*, 596.
- [4] S. K. Deb, *Sol. Energy Mater. Sol. Cells* **1992**, *250*, 327.
- [5] R. S. Bonilla, B. Hoex, P. Hamer, P. R. Wilshaw, *Phys. Status Solidi A* **2017**, *2140*, 1700293.
- [6] G. Dingemans, F. Einsele, W. Beyer, M. C. M. van de Sanden, W. M. M. Kessels, *J. Appl. Phys.* **2012**, *1110*, 093713.
- [7] B. Hoex, J. Schmidt, P. Pohl, M. C. M. van de Sanden, W. M. M. Kessels, *J. Appl. Phys.* **2008**, *1040*, 0044903.
- [8] S. Huang, Q. Wu, Z. Jia, X. Jin, X. Fu, H. Huang, X. Zhang, J. Yao, J. Xu, *Adv. Opt. Mater.* **2020**, *80*, 1901808.
- [9] A. S. Mayet, H. Cansizoglu, Y. Gao, S. Ghandiparsi, A. Kaya, C. Bartolo-Perez, B. AlHalaili, T. Yamada, E. P. Devine, A. F. Elrefaie, S.-Y. Wang, M. S. Islam, *J. Opt. Soc. Am. B* **2018**, *350*, 1059.
- [10] H. J. Queisser, E. E. Haller, *Science* **1998**, *2810*, 945.
- [11] P. Repo, J. Benick, V. Vähänissi, J. Schön, G. von Gastrow, B. Steinhäuser, M. C. Schubert, M. Hermle, H. Savin, *Energy Proc.* **2013**, *38*, 866.
- [12] P. Saint-Cast, Y.-H. Heo, E. Billot, P. Olwal, M. Hofmann, J. Rentsch, S. W. Glunz, R. Preu, *Energy Proc.* **2011**, *8*, 642.
- [13] S. Olibet, E. Vallat-Sauvain, C. Ballif, *Phys. Rev. B* **2007**, *760*, 035326.
- [14] M. J. Kerr, A. Cuevas, *Semicond. Sci. Technol.* **2001**, *170*, 35.
- [15] B. Hoex, J. J. H. Gielis, M. C. M. van de Sanden, W. M. M. Kessels, *J. Appl. Phys.* **2008**, *1040*, 113703.
- [16] V. Naumann, M. Otto, R. B. Wehrspohn, M. Werner, C. Hagendorf, *Energy Proc.* **2012**, *27*, 312.

- [17] F. Werner, in *Atomic Layer Deposition of Aluminum Oxide on Crystalline Silicon: Fundamental Interface Properties and Application to Solar Cells*, **2014**, ISBN 978-3-8440-3039-6.
- [18] B. W. Busch, O. Pluchery, Y. J. Chabal, D. A. Muller, R. L. Opila, J. R. Kwo, E. Garfunkel, *MRS Bull.* **2002**, 270, 206.
- [19] M. N. Getz, M. Povoli, E. Monakhov, *J. Appl. Phys.* **2021**, 1290, 205701.
- [20] F. Werner, B. Veith, D. Zielke, L. Kühnemund, C. Tegenkamp, M. Seibt, R. Brendel, J. Schmidt, *J. Appl. Phys.* **2011**, 1090, 113701.
- [21] N. M. Terlinden, G. Dingemans, M. C. M. van de Sanden, W. M. M. Kessels, *Appl. Phys. Lett.* **2010**, 960, 112101.
- [22] D. K. Simon, P. M. Jordan, T. Mikolajick, I. Dirnstorfer, *ACS Appl. Mater. Interfaces* **2015**, 70, 28215.
- [23] M. Garin, J. Heinonen, L. Werner, T. P. Pasanen, V. Vähänissi, A. Haarahiltunen, M. A. Juntunen, H. Savin, *Phys. Rev. Lett.* **2020**, 1250, 117702.
- [24] E. Grusell, S. Berg, L. P. Andersson, *J. Electrochem. Soc.* **1980**, 1270, 1573.
- [25] L. P. Andersson, A. O. Ewaraye, *Vacuum* **1978**, 280, 5.
- [26] O. Paz, F. D. Aurret, J. F. White, *J. Electrochem. Soc.* **1984**, 1310, 1712.
- [27] F. Volpi, A. R. Peaker, I. Berbezier, A. Ronda, *J. Appl. Phys.* **2004**, 950, 4752.
- [28] F. D. Aurret, P. M. Mooney, *J. Appl. Phys.* **1984**, 550, 988.
- [29] B. El-Kareh, A. F. Puttlitz, *J. Vac. Sci. Technol.* **1978**, 150, 1047.
- [30] K. F. Galloway, S. Mayo, *Solid State Technol.* **1979**, 22, 96.
- [31] D. V. Lang, *J. Appl. Phys.* **1974**, 450, 3023.
- [32] K. A. Collett, R. S. Bonilla, P. Hamer, G. Bourret-Sicotte, R. Lobo, T. Kho, P. R. Wilshaw, *Sol. Energy Mater. Sol. Cells* **2017**, 173, 50.
- [33] K. A. Collett, M. Cyrson, R. S. B. Osorio, P. R. Wilshaw, in *European Photovoltaic Solar Energy Conf. and Exhibition* **2016**.
- [34] B. E. Deal, *Silicon Dev. Process. Proc.* **1969**, 13, 36.
- [35] W. D. Eades, *Ph.D. Thesis*, Stanford University, **1985**.
- [36] O. E. Setälä, T. P. Pasanen, J. Ott, V. Vähänissi, H. Savin, *Phys. Status Solidi A* **2021**, 2180, 2100214.
- [37] A. C. Diebold, in *Handbook of Silicon Semiconductor Metrology*, CRC Press, Boca Raton, FL **2001**.
- [38] L. M. Terman, *Solid-State Electron.* **1962**, 50, 285.
- [39] M. Wilson, J. Lagowski, L. Jastrzebski, A. Savtchouk, V. Faifer, *AIP Conf. Proc.* **2001**, 5500, 220.
- [40] D. K. Schroder, *Mater. Sci. Eng. B* **2002**, 91, 196.
- [41] A. M. Hoff, D. K. DeBusk, R. W. Schanzer, in *In-Line Methods and Monitors for Process and Yield Improvement*, Vol 3884, International Society for Optics and Photonics, Bellingham, WA **1999**, pp. 207–215.
- [42] D. K. Schroder, in *Semiconductor Material and Device Characterization*, John Wiley & Sons, Inc., Hoboken, NJ **2006**.
- [43] B. E. Deal, E. L. MacKenna, P. L. Castro, *J. Electrochem. Soc.* **1969**, 1160, 997.
- [44] R. Khosla, D. Schwarz, H. S. Funk, K. Guguieva, J. Schulze, *Solid-State Electron.* **2021**, 185, 108027.
- [45] N. M. Johnson, D. K. Biegelsen, M. D. Moyer, *J. Vac. Sci. Technol.* **1981**, 190, 390.
- [46] G. Seguini, E. Ciani, C. Wiemer, D. Saynova, J. A. M. van Roosmalen, M. Perego, *Appl. Phys. Lett.* **2013**, 102, 131603.
- [47] K.-H. Cho, Y. J. Cho, H. S. Chang, K.-J. Kim, H. E. Song, *J. Korean Phys. Soc.* **2015**, 670, 995.
- [48] I. Ionica, D. Damianos, A. Kaminski, D. Blanc-Péllissier, G. Ghibaudo, S. Cristoloveanu, L. Bastard, A. Bouchard, X. Mescot, M. Gri, M. Lei, B. Larzelere, G. Vitrant, *ECS Trans.* **2020**, 970, 119.