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## A posteriori analysis of classical plate elements

Tom Gustafsson, Rolf Stenberg<sup>1</sup> and Juha Videman

**Summary.** We outline the results of our recent article on the a posteriori error analysis of  $C^1$  finite elements for the classical Kirchhoff plate model with general boundary conditions. Numerical examples are given.

*Key words:* Kirchhoff plate model,  $C^1$  elements, a posteriori error estimates.

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### Introduction

The purpose of our work is to fill a gap in the literature. Surprisingly, the a posteriori error analysis for classical plate finite elements has so far only been given for the fully clamped case and a load in  $L^2$ , cf. [2]. In our recent work [1], we treated a combination of all common boundary conditions (clamped, simply supported and free). In addition, we considered the cases of point and line loads.

### The Kirchhoff plate problem

We denote the deflection of the plate's midsurface by  $u$ , the curvature by  $\mathbf{K}$  and the moment by  $\mathbf{M}$ , and we assume isotropic linear elasticity. Hence, it holds

$$\mathbf{M}(u) = \frac{d^3}{12} \mathbb{C} \mathbf{K}(u), \quad (1)$$

with

$$\mathbb{C} \mathbf{A} = \frac{E}{1 + \nu} \left( \mathbf{A} + \frac{\nu}{1 - \nu} (\text{tr } \mathbf{A}) \mathbf{I} \right), \quad \forall \mathbf{A} \in \mathbb{R}^{2 \times 2}, \quad (2)$$

where  $d$  denotes the thickness of the plate.  $E$  and  $\nu$  are the Young's modulus and Poisson ratio, respectively. The strain energy for an admissible deflection  $v$  is then  $\frac{1}{2}a(v, v)$ , with

$$a(w, v) = \int_{\Omega} \mathbf{M}(w) : \mathbf{K}(v) \, dx = \int_{\Omega} \frac{d^3}{12} \mathbb{C} \boldsymbol{\varepsilon}(\nabla w) : \boldsymbol{\varepsilon}(\nabla v) \, dx. \quad (3)$$

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The potential energy  $l(v)$  stems from the loading, which we assume to consist of a distributed load  $f \in L^2(\Omega)$ , a load  $g \in L^2(S)$  along the line  $S \subset \Omega$ , and of a point load  $F$  at the point  $x_0$ , so that

$$l(v) = \int_{\Omega} f v \, dx + \int_S g v \, ds + F v(x_0). \quad (4)$$

The total energy is thus  $\frac{1}{2}a(v, v) - l(v)$ , and its minimisation leads to the variational form: *find*  $u \in V$  *such that*

$$a(u, v) = l(v) \quad \forall v \in V, \quad (5)$$

with

$$V = \{v \in H^2(\Omega) \mid v|_{\Gamma_c \cup \Gamma_s} = 0, \quad \frac{\partial u}{\partial n}|_{\Gamma_c} = 0\}. \quad (6)$$

We assume that the plate is clamped on the boundary part  $\Gamma_c$ , simply supported on  $\Gamma_s$ , and free on  $\Gamma_f = \partial\Omega \setminus (\Gamma_c \cup \Gamma_s)$ .

By the well-known integration by parts, we get the boundary value problem. To this end we have to recall the following quantities for a admissible displacement  $v$ ; the normal shear force  $Q_n(v)$ , the normal and twisting moments  $M_{nn}(v)$ ,  $M_{ns}(v)$ , and the effective shear force

$$V_n(v) = Q_n(v) + \frac{\partial M_{ns}(v)}{\partial s}. \quad (7)$$

With the constitutive relationship (2), an elimination yields the plate equation for the deflection  $u$ :

$$\mathcal{A}(u) := D\Delta^2 u = l, \quad (8)$$

where the so-called bending stiffness  $D$  is defined as

$$D = \frac{Ed^3}{12(1-\nu^2)}. \quad (9)$$

The boundary value problem is the following.

- In the domain we have the distributional *differential equation*

$$\mathcal{A}(u) = l \quad \text{in } \Omega, \quad (10)$$

where  $l$  is the distribution defined by (4).

- On the *clamped* part we have the conditions:  $u = 0$  and  $\frac{\partial u}{\partial n} = 0$  on  $\Gamma_c$ .
- On the *simply supported* part it holds:  $u = 0$  and  $M_{nn}(u) = 0$  on  $\Gamma_s$ .
- On the *free part* it holds:  $M_{nn}(u) = 0$  and  $V_n(u) = 0$  on  $\Gamma_f$ .
- At the *corners on the free part* we have the jump condition on the twisting moment

$$[[M_{ns}(u)(c)]] = 0 \quad \text{for all corners } c \in \Gamma_f.$$

Here and below  $[[\cdot]]$  denotes the jump.

We consider conforming finite element methods: *find*  $u_h \in V_h \subset V$  *such that*

$$a(u_h, v) = l(v) \quad \forall v \in V_h. \quad (11)$$

The finite element partitioning is denoted by  $\mathcal{C}_h$ . We assume that mesh is such that the point load is a vertex and the line load is along edges. The edges are divided into interior edges  $\mathcal{E}_h^i$ , edges on  $S$ ,  $\mathcal{E}_h^S$ , edges on the free boundary  $\mathcal{E}_h^f$ , and edges on the simply supported boundary  $\mathcal{E}_h^s$ . The local error indicators are then the following.

- The residual on each element

$$h_K^2 \|\mathcal{A}(u_h) - f\|_{0,K}, \quad K \in \mathcal{C}_h.$$

- The jump residuals of the normal moment along interior edges

$$h_E^{1/2} \|[[M_{nn}(u_h)]]\|_{0,E}, \quad E \in \mathcal{E}_h^i.$$

- The jump residuals in the effective shear force along interior edges

$$h_E^{3/2} \|[[V_n(u_h)]] - g\|_{0,E}, \quad E \in \mathcal{E}_h^S, \quad h_E^{3/2} \|[[V_n(u_h)]]\|_{0,E}, \quad E \in \mathcal{E}_h^i \setminus \mathcal{E}_h^S.$$

- The normal moment along edges on the free and simply supported boundaries

$$h_E^{1/2} \|M_{nn}(u_h)\|_{0,E}, \quad E \in \mathcal{E}_h^f \cup \mathcal{E}_h^s.$$

- The effective shear force along edges on the free boundary

$$h_E^{3/2} \|V_n(u_h)\|_{0,E}, \quad E \in \mathcal{E}_h^f.$$

The error estimator is defined through

$$\begin{aligned} \eta^2 = & \sum_{K \in \mathcal{C}_h} h_K^4 \|\mathcal{A}(u_h) - f\|_{0,K}^2 + \sum_{E \in \mathcal{E}_h^S} h_E^3 \|[[V_n(u_h)]] - g\|_{0,E}^2 + \sum_{E \in \mathcal{E}_h^i \setminus \mathcal{E}_h^S} h_E^3 \|[[V_n(u_h)]]\|_{0,E}^2 \\ & + \sum_{E \in \mathcal{E}_h^i} h_E \|[[M_{nn}(u_h)]]\|_{0,E}^2 + \sum_{E \in \mathcal{E}_h^f} h_E^3 \|V_n(u_h)\|_{0,E}^2 + \sum_{E \in \mathcal{E}_h^f \cup \mathcal{E}_h^s} h_E \|M_{nn}(u_h)\|_{0,E}^2. \end{aligned} \quad (12)$$

Our a posteriori estimate is the following, where the energy norm is defined as  $\|v\| = a(v, v)^{1/2}$ .

**Theorem 1** *There exists positive constants  $C_1, C_2$ , such that*

$$C_1 \eta \leq \|u - u_h\| \leq C_2 \eta. \quad (13)$$

## Numerical examples

In the examples, we have used the Argyris triangle. In the figures, we give the meshes for the adaptive solution of a square plate with a point and line load, and for a L-shaped domain with a free boundary for the edges sharing the re-entrant corner and simply supported along the rest of the boundary.

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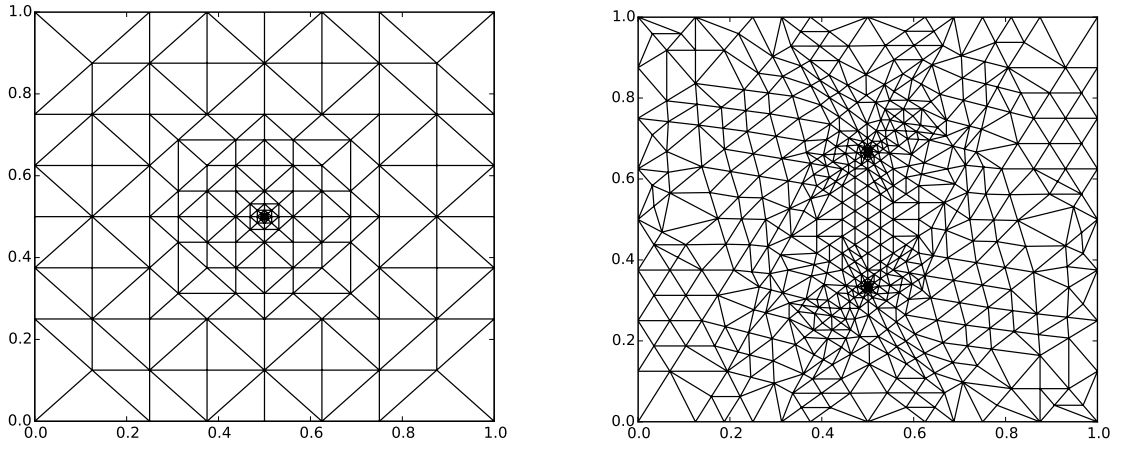


Figure 1. The adaptive meshes for the point and line loads.

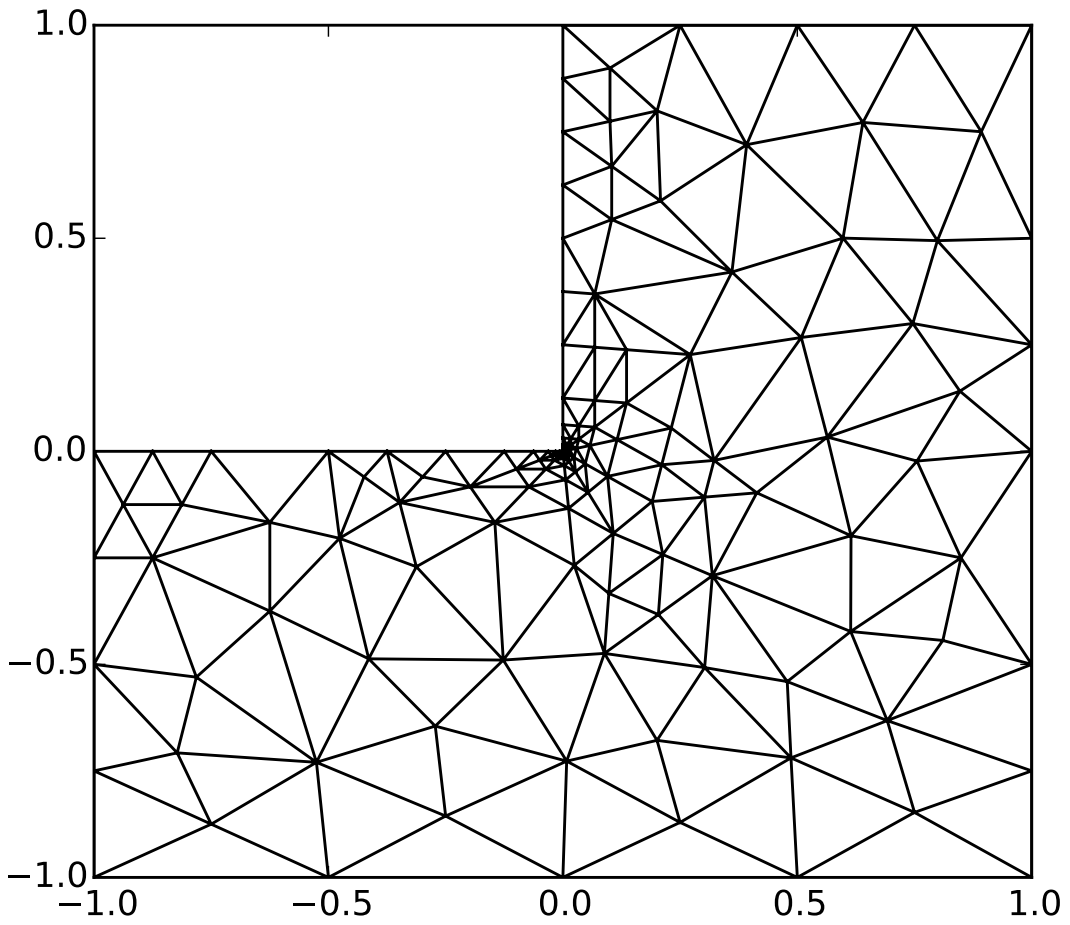


Figure 2. The adaptive mesh for the L-shaped domain.

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