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Published in:
Physical Review B

DOI:
[10.1103/PhysRevB.93.064403](https://doi.org/10.1103/PhysRevB.93.064403)

Published: 02/02/2016

Document Version
Publisher's PDF, also known as Version of record

Please cite the original version:
Estevez , V., & Laurson, L. (2016). Magnetic domain-wall dynamics in wide permalloy strips. *Physical Review B*, 93(6), Article 064403. <https://doi.org/10.1103/PhysRevB.93.064403>

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Magnetic domain-wall dynamics in wide permalloy strips

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(Received 18 November 2015; published 2 February 2016)

Domain walls in soft permalloy strips may exhibit various equilibrium micromagnetic structures depending on the width and thickness of the strip, ranging from the well-known transverse and vortex walls in narrow and thin strips to double and triple vortex walls recently reported in wider strips [V. Estévez and L. Laurson, *Phys. Rev. B* **91**, 054407 (2015)]. Here, we analyze the field driven dynamics of such domain walls in permalloy strips of widths from 240 nm up to 6 μm , using the known equilibrium domain wall structures as initial configurations. Our micromagnetic simulations show that the domain wall dynamics in wide strips is very complex, and depends strongly on the geometry of the system, as well as on the magnitude of the driving field. We discuss in detail the rich variety of the dynamical behaviors found, including dynamic transitions between different domain wall structures, periodic dynamics of a vortex core close to the strip edge, transitions towards simpler domain wall structures of the multi-vortex domain walls controlled by vortex polarity, and the fact that for some combinations of the strip geometry and the driving field the system cannot support a compact domain wall.

DOI: 10.1103/PhysRevB.93.064403

I. INTRODUCTION

The study of static and dynamic properties of the magnetic domain walls (DWs) in ferromagnetic nanostructures has recently attracted a lot of attention not only due to the related fundamental physics aspects, but also because of the emergence of technological applications based on DWs and their dynamics, such as memory [1,2] and logic devices [3–5]. Consequently, DW dynamics in nanostrips and wires has been extensively studied, using both applied magnetic fields [6–15] and spin-polarized electric currents [16–21] as the driving force.

One of the typical materials used for the study of DWs and their dynamics in strip geometry is permalloy, a soft magnetic material with a negligible magnetocrystalline anisotropy [23,24]. Such systems exhibit in-plane domains along the long axis of the strip, induced by shape anisotropy. Previous works have shown that the structures of the DWs separating such domains may be relatively complex, and depend strongly on the sample geometry [9,10,25–27]. For nanostrip geometries, with the strip width and thickness of the order of 100 nm and 10 nm or less, respectively, the known equilibrium DW structures are the transverse DW (TW) and the asymmetric transverse DW (ATW) [27–33]. For somewhat wider and/or thicker strips, the vortex DW (VW) is the stable structure [27,30–34]. A related, three-dimensional vortexlike flux-closing DW structure has been reported in permalloy strips of thickness exceeding 60 nm [35]. Recently, we have shown that in the case of even wider permalloy strips than those with VW as the equilibrium DW structure, two additional equilibrium DW structures appear, namely the double vortex (DVW) and triple vortex walls (TVW) [26].

Studies of field and current driven dynamics of DWs in permalloy strips have focused mostly on narrow strips with either TW or VW as the equilibrium DW structure [8–13,36–41]. In both cases, the DW dynamics exhibits a Walker breakdown [42], an instability occurring when the DW

internal degrees of freedom are excited by a strong enough external driving force. This force can be a magnetic field $B > B_W$ or a spin-polarized electric current $J > J_W$, with B_W and J_W the Walker field and current density, respectively [8,11,15,22]. For small driving forces, the DW velocity increases up to B_W or J_W (steady, or viscous regime). At the onset of the Walker breakdown, the DW velocity decreases abruptly as a consequence of periodic transitions between different DW structures. In narrow and thin strips with TW as the equilibrium structure, repeated transitions between TWs of different polarities (or signs of the TW internal magnetization) take place via nucleation and propagation of an antivortex across the strip width; we will refer to the transient state as the antivortex wall, or AVW [10]. For wider strips with a VW equilibrium DW, periodic transitions between VW and TW structures are typically observed [10,36]. However, so far, little is known about the DW dynamics in even wider strips, especially ones with DVW or TVW as the equilibrium DW structure.

In this paper, we analyze numerically field driven DW dynamics in permalloy strips considering a wide range of strip widths w from 240 nm up to 6 μm . We focus on relatively thin strips with the thickness Δ_z in the range of 5 to 25 nm. Depending on w and Δ_z , the equilibrium DW structure in such strips can be either VW, DVW, or TVW [26] (see Fig. 1 for examples of these structures). Since the space of parameters spanned by w , Δ_z , and B_{ext} is quite large, we focus on specific example cases, covering each of the three above-mentioned equilibrium DW structures, and illustrating the complex nature of the DW dynamics in wide strips, exhibiting also a strong dependence on both the geometry of the system, and on the magnitude of the driving field. Already for rather confined geometries with VW as the equilibrium DW structure, several possible dynamical behaviors are encountered, in particular for large enough B_{ext} (i.e., above the small field steady/viscous regime). In addition to various periodic transitions occurring between different DW structures, we also report on localized oscillatory motion of the vortex core close to the strip edge. This novel behavior is characteristic of wide enough strips, giving rise to an extended plateau in the velocity versus field

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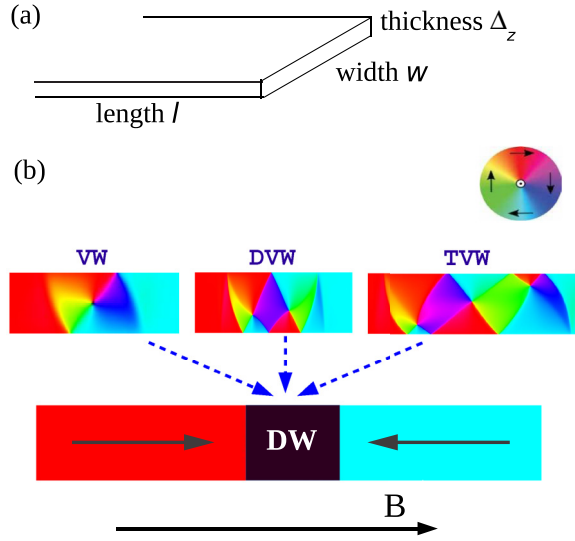


FIG. 1. (a) Geometry of the permalloy strip. (b) A top view of the magnetization in the initial state. Magnetization points along the long axis of the strip within the two domains (as indicated by the arrows) forming a head-to-head configuration. The domain wall separating these domains has a geometry-dependent equilibrium structure. Here, we consider systems where the latter is either a vortex wall (VW), double vortex wall (DVW), or a triple vortex wall (TVW). To study the DW dynamics, an external magnetic field B_{ext} is applied along the long axis of the strip.

curve. Another signature of DW dynamics in wide strips is that for some combinations of w , Δ_z , and B_{ext} , DWs are unstable: the DW's width grows without limit as different components of the DW assume a different propagation velocity [14]. DVWs are found to be structurally stable against small applied fields depending on the polarity of the vortices. In the stable case (parallel polarity), the dynamics exhibits a unique small field steady/viscous regime, but larger fields tend to induce a transition into a simpler VW structure, or lead to complex, turbulentlike dynamics with repeated transitions between numerous different DW structures. For antiparallel polarity, the DVW transforms in a VW even for very small fields. For larger fields the dynamics observed is the same independently of the polarity. The same is true also for TVWs, whose stability at low fields depends on the polarities of the three vortices. In general, the TVWs tend to transform into simpler DW structures during their field-driven dynamics. The paper is organized as follows. The following section (Sec. II) describes the details of our micromagnetic simulations, while in the three sections after that (Secs. III, IV, and V), DW dynamics starting from VW, DVW, and TVW equilibrium initial configurations, respectively, are described and discussed. Section VI finishes the paper with a summary and conclusions.

II. MICROMAGNETIC SIMULATIONS

The system under study consists of a permalloy strip of width w and thickness Δ_z , satisfying $\Delta_z \ll w$, see Fig. 1(a). To mimic an infinitely long strip, the magnetic charges at the two ends of the strip are compensated. In all the cases

considered the actual simulated length l , i.e., the length of the computational window which is centered at and moving with the DW, satisfies $l \geq 4w$. The initial state is an in-plane head-to-head domain structure, with the equilibrium DW structure [26] corresponding to the w and Δ_z of the strip in the middle of the sample, see Fig. 1(b). To analyze the DW dynamics, a constant external magnetic field B_{ext} is applied along the long axis of the strip. All the results presented in this work have been obtained using the typical material parameters of permalloy, i.e., saturation magnetization $M_s = 860 \times 10^3$ A/m, exchange constant $A_{\text{ex}} = 13 \times 10^{-12}$ J/m, and the Gilbert damping constant $\alpha = 0.01$. We focus on the ideal case of perfect strips without any kind of disorder or impurities, and consider a temperature T equal to zero.

The micromagnetic simulations have been performed using the GPU-based micromagnetic code MUMAX3 [43–45]. To calculate the magnetization dynamics of the system, the Landau-Lifshitz-Gilbert equation [46,47],

$$\partial \mathbf{m} / \partial t = \gamma \mathbf{H}_{\text{eff}} \times \mathbf{m} + \alpha \mathbf{m} \times \partial \mathbf{m} / \partial t, \quad (1)$$

is solved numerically. Here, \mathbf{m} is the magnetization, γ the gyromagnetic ratio, and \mathbf{H}_{eff} the effective field, with contributions due to exchange, Zeeman, and demagnetizing energies. The size of the discretization cell used depends on the system size, ranging from 3 to 5 nm, but is always bounded by the exchange length, $\Lambda = (2A/\mu_0 M_s^2)^{1/2} \approx 5$ nm, in the in-plane directions, and equals Δ_z in the out-of-plane direction. The DW velocity is calculated once the DW structure is in the steady state; thus, the total simulation time considered is at least 300 ns.

III. VORTEX WALL DYNAMICS

The vortex wall is a DW structure characterized by the chirality or sense of rotation of the vortex, and the core polarity p . The latter may assume two different values, $p = \pm 1$, corresponding to the two possible out-of-plane magnetization directions of the core. Vortex wall is the equilibrium DW structure within a wide range of w and Δ_z [26], and thus it is a pertinent question to what extent the dynamics exhibited by it depends on the geometry of the strip. To this end, in what follows, we consider some example cases, demonstrating that already in relatively narrow strips, several different kinds of field-driven VW dynamics may be observed.

A. Vortex wall dynamics in narrow strips

To illustrate this, we start by considering the DW velocity v as a function of the applied field B_{ext} . In Fig. 2, we show that quite different $v(B_{\text{ext}})$ curves are obtained depending on w and Δ_z . Figure 2(a) shows that for a relatively narrow strip width $w = 240$ nm, considering two different thicknesses ($\Delta_z = 10$ and 20 nm) leads to clear differences in the $v(B_{\text{ext}})$ curves already below the Δ_z dependent Walker field. In the case of $\Delta_z = 10$ nm, for very small fields, v increases linearly with B_{ext} , a behavior arising from steady VW motion with the vortex core assuming an off-center position within the strip [10]. For $B_{\text{ext}} \geq 0.4$ mT, the vortex core is expelled out of the strip due to the gyrotropic force [48], leading to a TW, and to a sublinear $v(B_{\text{ext}})$ relation [10]. In contrast, for the thicker strip with

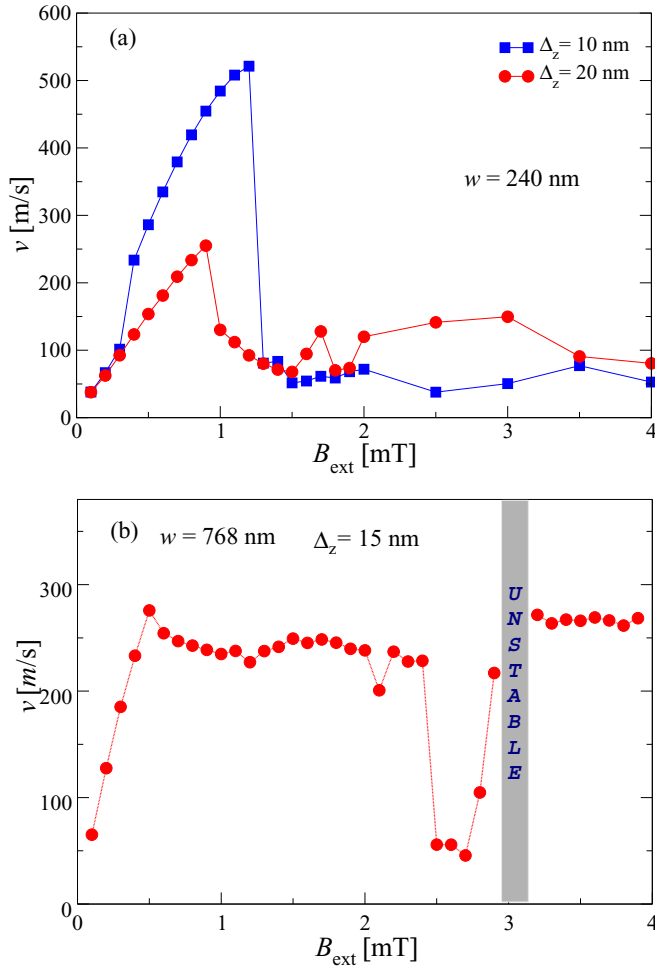


FIG. 2. $v(B_{\text{ext}})$ curves for different strip geometries, where the VW is the equilibrium DW structure. (a) For two strips, both having the same width of $w = 240$ nm and in-plane cell size 5 nm, and different thicknesses $\Delta_z = 10, 20$ nm. (b) The same for a strip with $w = 768$ nm, $\Delta_z = 15$ nm and in-plane cell size 3 nm. For $3 \text{ mT} \leq B_{\text{ext}} \leq 3.1 \text{ mT}$ the DW is unstable; the DW width grows without bound.

$\Delta_z = 20$ nm, the linear dependence of v on B_{ext} , originating from steady VW dynamics with the vortex core within the strip, persists up to the Walker field. Notice also that for B_{ext} such that the VW has transformed into a TW for $\Delta_z = 10$ nm, the TW velocity for a given B_{ext} exceeds the corresponding VW velocity in the system with $\Delta_z = 20$ nm. This is likely due to the large energy dissipation associated with the dynamics of the vortex core [48].

Above the Walker breakdown, v drops dramatically as a consequence of the onset of periodic transformations between different types of DW structures. An example of this behavior is shown in Fig. 3, where the VW dynamics for $B_{\text{ext}} = 1.5$ mT, exceeding the Walker field B_W , is shown for a strip with $w = 240$ nm and $\Delta_z = 10$ nm (see also Ref. [49], movie 1). Periodic transformations between VW and ATW structures are observed, mediated by the transverse motion of the vortex core across the strip width. Notice that here the vortex core dynamics alternates between injection from the top and bottom edges, followed by its top-to-bottom and bottom-to-top

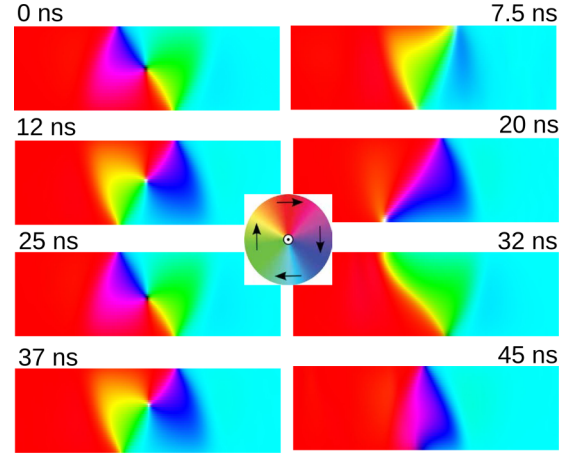


FIG. 3. Vortex wall dynamics in a strip with $w = 240$ nm and $\Delta_z = 10$ nm for $B_{\text{ext}} = 1.5$ mT (i.e., above the Walker field $B_W = 1.2$ mT). The DW dynamics repeats a cycle where the vortex core first moves out of the strip through the top strip edge, forming an ATW, followed by the injection of a vortex core from the top edge with the opposite polarity, which subsequently moves across the strip to the bottom edge, reversing the ATW magnetization. Then a vortex, again with a polarity opposite to that of the previous one, is nucleated from the bottom edge, and moves to the top edge, again reversing the ATW magnetization. Then the same process repeats.

propagation across the strip, respectively. In Ref. [10], the same geometry and B_{ext} were found to always lead to vortex core injection from the top strip edge, followed by its “downwards” motion towards the bottom edge. We have checked that this is due to slightly different values of M_s and A_{ex} (i.e., $M_s = 8 \times 10^5$ A/m and $A_{\text{ex}} = 10^{-11}$ J/m, respectively) used in the simulations of Ref. [10]. Given such sensitivity to small differences in the simulation conditions, one may expect that also other kinds of VW dynamics may be observed.

An example thereof is provided in Fig. 4, where the same geometry as above (i.e., $w = 240$ nm and $\Delta_z = 10$ nm) is considered for a larger driving field of $B_{\text{ext}} = 4$ mT. Now, the DW exhibits transformations between VW, ATW, and AVW structures, with the repeated ATW magnetization reversal mediated by alternating nucleation and propagation across the strip of either a vortex or an antivortex core (see also Ref. [49], movie 2). Thicker strips appear to lead to a smaller number of different periodicities. For instance, for $w = 240$ nm and $\Delta_z = 20$ nm, only two periodicities are observed for $B_{\text{ext}} > B_W$, consisting of transformations between VW and ATW structures. While the ATW reverses its magnetization, the VW structures exhibit alternating polarities but no change of chirality takes place (not shown). For some fields, also a nucleation of an antivortex in the edge of the strip is observed, followed by a vortex-antivortex annihilation process. Thus, already for relatively confined geometries, the field driven VW dynamics is quite complex, given that different geometries and applied fields may lead to numerous different dynamical behaviors.

B. Vortex wall dynamics in wide strips

In the case of wider strips, the VW dynamics is even more complex. Figure 2(b) shows the $v(B_{\text{ext}})$ curve for a strip with

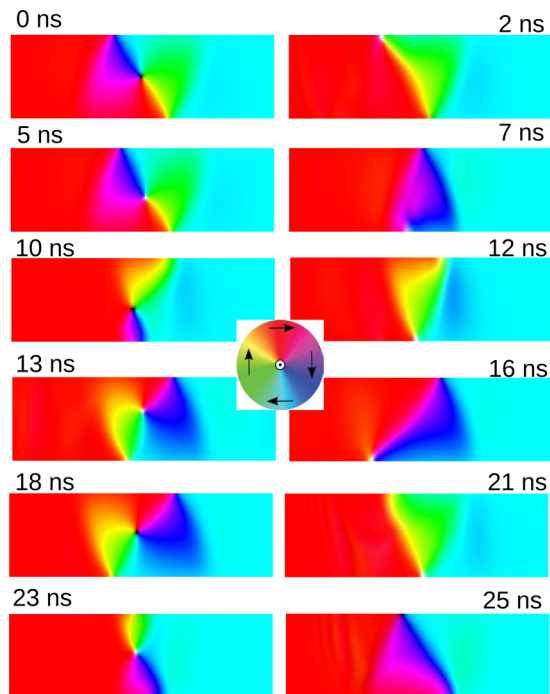


FIG. 4. Vortex wall dynamics in a strip with $w = 240$ nm and $\Delta_z = 10$ nm for $B_{\text{ext}} = 4$ mT (above the Walker breakdown, $B_W = 1.2$ mT). Now the dynamics alternates between ATWs of different magnetization, with its reversal mediated by either a vortex or an antivortex core moving across the strip width.

$w = 768$ nm and $\Delta_z = 15$ nm. As can be seen, the resulting $v(B_{\text{ext}})$ curve is very different to those observed for more narrow strips [Fig. 2(a)], and consists of six different regimes. Analogously to more narrow strips, at low fields the velocity increases linearly with B_{ext} , with the vortex core assuming again a steady, B_{ext} dependent off-center position.

For $0.6 \text{ mT} \leq B_{\text{ext}} \leq 2.4 \text{ mT}$, the $v(B_{\text{ext}})$ curve initially exhibits a small, negative gradient, but looks otherwise essentially like a plateau. Thus the DW is able to maintain a relatively high velocity of approximately 250 m/s, while for more narrow strips [Fig. 2(a)] Walker breakdown leads to significantly smaller DW velocities for most of this range of B_{ext} . This behavior arises as a consequence of the vortex core moving towards the strip edge due to the gyrotropic force, but contrary to the behavior observed in more narrow strips, it is not able to leave the strip. Instead, the competition between the gyrotropic force and a repulsive interaction of the vortex core with the half-antivortex (HAV) edge defect [14] leads to periodic oscillations of the vortex core y position in the vicinity of the strip edge. Figure 5 shows an example of the trajectory (y coordinate y_c as a function of time t) of the vortex core in a strip with $w = 768$ nm and $\Delta_z = 15$ nm with the DW subject to an applied field of $B_{\text{ext}} = 1.5$ mT. A snapshot of the DW configuration with the vortex core close to the HAV edge defect is included. An example of this process is also shown in Ref. [49], movie 3. The apparent emissions of spin waves observed in the movie may be related to short-lived (faster than the frame rate) transient dynamics involving nucleation/annihilation processes. Such behavior arises in wide strips due to the interplay between the gyrotropic

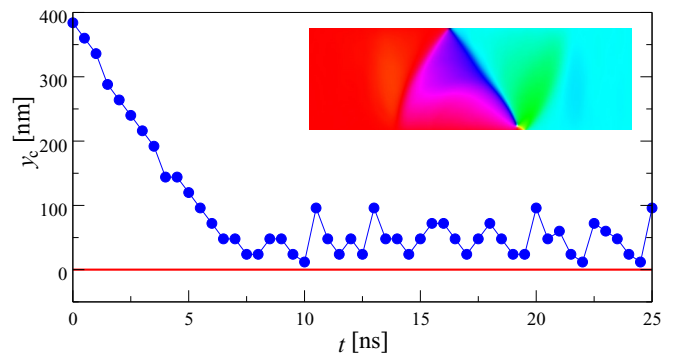


FIG. 5. The trajectory y_c of the vortex core in the vicinity of the edge of the system for $w = 768$ nm, $\Delta_z = 15$ nm, and $B_{\text{ext}} = 1.5$ mT, illustrating the attraction-repulsion effect. The red line at $y_c = 0$ corresponds to the strip edge. The inset shows an example of the DW configuration with the vortex core close to the edge of the strip.

force and the interaction of the vortex core with the HAV being different to the one in narrow strips. Since in permalloy everything is dominated by shape anisotropy originating from the strip edges, the energy needed to displace the vortex core from the middle of the strip towards the edge by the gyrotropic force is smaller than in narrow strips. This is evidenced also by the steep increase of $v(B_{\text{ext}})$ in Fig. 2(b) for small B_{ext} . At the same time, the HAV edge defect is more localized in relative terms in wider strips, thus leading to a relatively strong short-range vortex core-HAV interaction, preventing the former from being expelled from the strip. In other words, the energy necessary to move the vortex core from its equilibrium position in the middle of the strip towards the strip edge is smaller than the energy needed to push the core out of the strip. As a result, the core exhibits oscillatory dynamics close to the strip edge, and the DW assumes a relatively high propagation velocity. In what follows, we shall refer to this phenomenon as the attraction-repulsion effect.

This behavior is characteristic of all field values within the range $0.6 \text{ mT} \leq B_{\text{ext}} \leq 2.4 \text{ mT}$, with the exception of $B_{\text{ext}} = 2.1$ mT. Here, v is a little smaller than elsewhere in the plateau region visible in Fig. 2(b). For that case, although most of the time the usual attraction-repulsion effect takes place, in rare instances the repulsion effect is strong enough to displace the vortex core all the way to the other edge of the strip, leading to a structural DW transformation (VW to ATW), and to a reduced DW propagation velocity.

For $2.5 \text{ mT} \leq B_{\text{ext}} \leq 2.7 \text{ mT}$, v decreases abruptly [Fig. 2(b)]. In this regime, the DW dynamics exhibits repeated transitions between a VW and a kind of stretched TW (or ATW) structure, see Fig. 6(a). The peculiar feature of this Walker breakdownlike dynamics is that while the vortex core is able to reach the top edge of the strip, transforming the DW momentarily into a TW-like structure, the vortex core never reaches the opposite (bottom) edge of the strip. Thus the transient TW-like structure always has the same magnetization. The resulting oscillatory dynamics of the vortex core resembles that observed above in the context of the attraction-repulsion effect. However, the amplitude of the transverse core oscillations is significantly larger, leading to a higher rate of energy dissipation associated with the internal

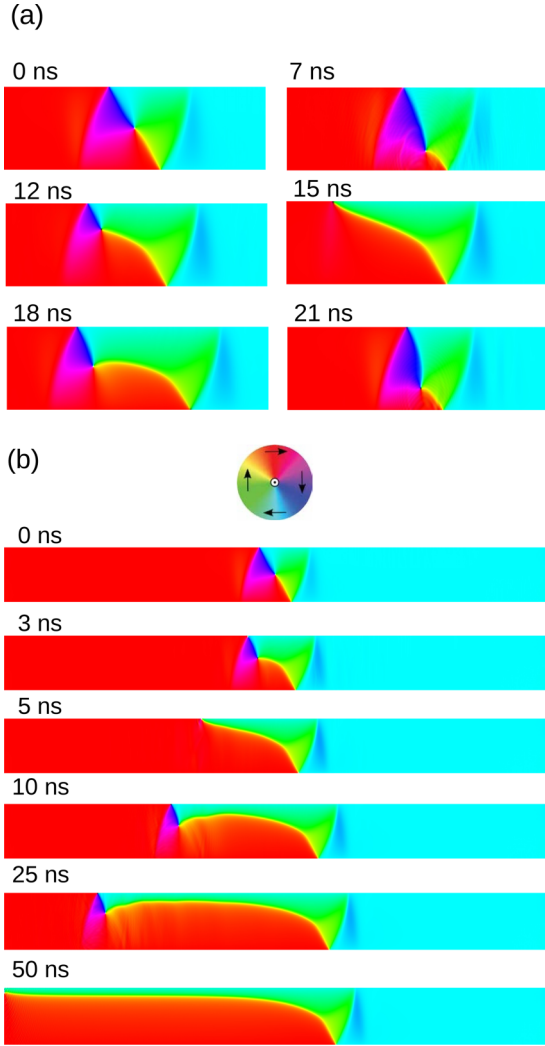


FIG. 6. (a) Walker breakdown-type of dynamics for $w = 768$ nm, $\Delta_z = 15$ nm, and $B_{\text{ext}} = 2.5$ mT. The vortex core exhibits asymmetric lateral oscillations, leaving the strip at the top edge, but reversing the direction of motion and polarity of the core well inside the strip when closer to the bottom edge. (b) The instability occurring for $B_{\text{ext}} = 3$ mT, where the DW width grows without limit as the leading HAV edge defect moves faster than the combination of the trailing HAV and the vortex core.

dynamics of the DW, resulting also in a significantly reduced DW propagation velocity.

For $B_{\text{ext}} = 2.8$ and 2.9 mT, the velocity increases as the attraction-repulsion effect appears again. In the case of $B_{\text{ext}} = 2.8$ mT, although also an attraction-repulsion effect is observed, the transformations between different DW structures (VW and ATW) are predominant. For $B_{\text{ext}} = 2.9$ mT, only the attraction-repulsion effect is observed, leading to a higher DW propagation velocity [Fig. 2(b)].

For $B_{\text{ext}} = 3$ and 3.1 mT, the leading and trailing edges of the DW start to move at different velocities, implying that the DW width grows without limit, and thus the system is unable to support a compact DW structure. The leading HAV moves faster than the trailing one, and the vortex core exhibits oscillatory transverse motion in the vicinity of the latter, see Fig. 6(b). Eventually the trailing edge of the DW

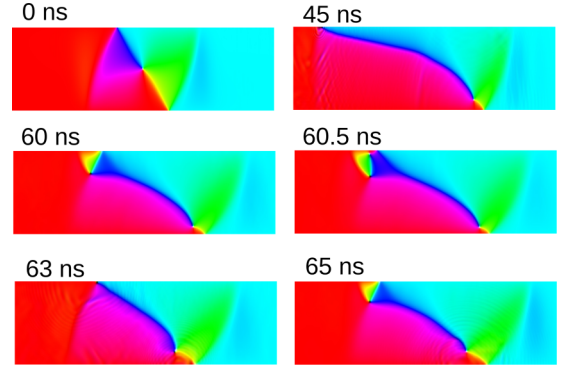


FIG. 7. VW dynamics for $B_{\text{ext}} = 3.3$ mT in a strip with $w = 768$ nm and $\Delta_z = 15$ nm. The trailing HAV edge defect repeatedly emits a vortex-antivortex pair, which then annihilates, and the process is repeated.

exits the computational window. Such unstable DW dynamics has previously been reported in Ref. [14], and renders the measurement of a well-defined DW velocity impossible.

Curiously, for larger fields ($3.2 \text{ mT} \leq B_{\text{ext}} \leq 4 \text{ mT}$; we focus here on fields up to 4 mT) the DW structure is stable again, and the $v(B_{\text{ext}})$ curve displays a second high-velocity plateau with v exceeding 250 m/s . There, after a relatively long initial transient, the DW dynamics proceeds as follows. In the vicinity of the leading edge HAV, the vortex core exhibits similar oscillatory dynamics as in the attraction-repulsion effect. The trailing HAV repeatedly emits a vortex-antivortex pair, which subsequently annihilates. Figure 7 and movie 4 in Ref. [49] illustrate this dynamics. Notice that while the above-mentioned sequence is the dominating one, sometimes the vortex-antivortex pair fails to annihilate in the vicinity of the trailing HAV, and is able to move to the opposite edge, leading to a short period of more complex dynamics before the above-mentioned sequence starts again (see Ref. [49], movie 4). Similar vortex-antivortex annihilation processes within DWs have previously been reported in strips with much more confined lateral dimensions, and applied fields roughly five times as high as here [50].

Vortex-antivortex annihilation processes can be parallel or antiparallel depending on the polarity of the vortex-antivortex pair [51–53]. When the vortex and the antivortex have the same polarity (parallel case), the process is continuous, whereas for the antiparallel one there is emission of spin waves as a result of the annihilation. This is related to the topological charges (winding number n) and the polarities of the defects. DW structures are composed of topological defects, which have associated winding numbers $n = +1$ for vortices, -1 for antivortices, and $\pm 1/2$ for edge defects [54]. In DW structures the topological defects are compensated, leading to zero total winding number, see Fig. 12. However, there are also other “topological charges” associated with the winding number and the polarity of the core of vortices and antivortices, e.g., the skyrmion charge $q = np/2$ [55]. For parallel polarity, vortex and antivortex have opposite skyrmion numbers, leading to a continuous annihilation process. However, for the antiparallel case the skyrmion numbers are equal, and there is no compensation. As a result, emission of spin waves occurs after the annihilation process [51–53]. Here, we have observed

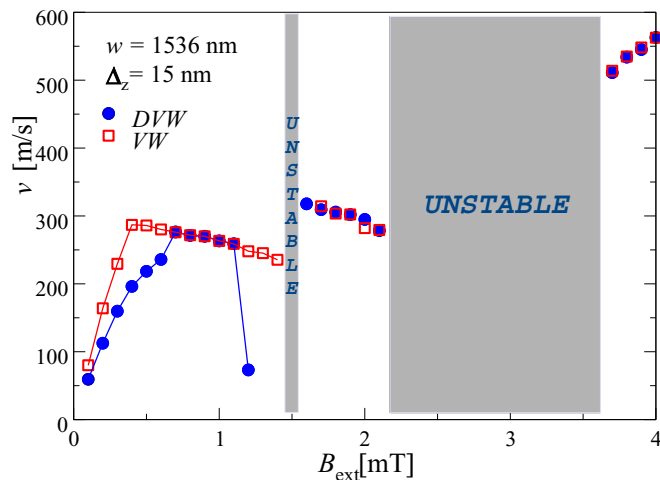


FIG. 8. $v(B_{\text{ext}})$ curves for a VW and a DVW initial state in a strip with $w = 1536$ nm and $\Delta_z = 15$ nm, where both structures are equilibrium DWs. In-plane cell size equal to 3 nm. The gray areas indicate unstable regions (no compact DWs).

both parallel and antiparallel processes of annihilation of the vortex-antivortex pair.

For even larger strips the VW dynamics is again found to be different from that observed in more narrow systems. For example, let us consider a strip with $w = 1536$ nm and $\Delta_z = 15$ nm, for which both VW and DVW are equilibrium structures (i.e., they have the same energy [26]). The $v(B_{\text{ext}})$ curve obtained when starting from a VW initial state is shown as red squares in Fig. 8. As before, the low-field steady/viscous regime (for $B_{\text{ext}} \leq 0.4$ mT) is followed by a plateaulike regime where the velocity decreases slowly with the field (for $0.5 \text{ mT} \leq B_{\text{ext}} \leq 1.4$ mT), due to the attraction-repulsion effect. Contrary to the case of the more narrow strip considered in Fig. 2(b), here no Walker breakdown is observed. Instead, the attraction-repulsion regime is terminated by an instability where the DW width grows without bound for $B_{\text{ext}} = 1.5$ and 1.6 mT.

Above this unstable regime, the velocity again decreases linearly with the field ($1.7 \text{ mT} \leq B_{\text{ext}} \leq 2.1$ mT). An example of the type of DW dynamics observed in this regime is shown in Fig. 9 for $B_{\text{ext}} = 2.1$ mT (see also Ref. [49], movie 5). When the field is applied the vortex core is displaced from its equilibrium position, and an antivortex enters into the strip. After that, a vortex nucleates into the strip. The antivortex moves towards the previous vortex, and an annihilation process between them takes place (snapshot at 34 ns shows the DW structure just after the annihilation). Unlike in the case observed for the strip with $w = 768$ nm and $\Delta_z = 15$ nm, where the processes of annihilation were between the vortex and antivortex that nucleate into the strip, now the annihilation is between the antivortex and the vortex that was inside the strip previously. Once the annihilation process occurs, this part of the DW structure starts to move towards the other vortex, reducing the width of the DW structure. Again two parts of the DW structure are moving with different velocities. The vortex that does not participate in the annihilation is feeling an attraction-repulsion effect. Although this is the main dynamical sequence, sometimes other more complex

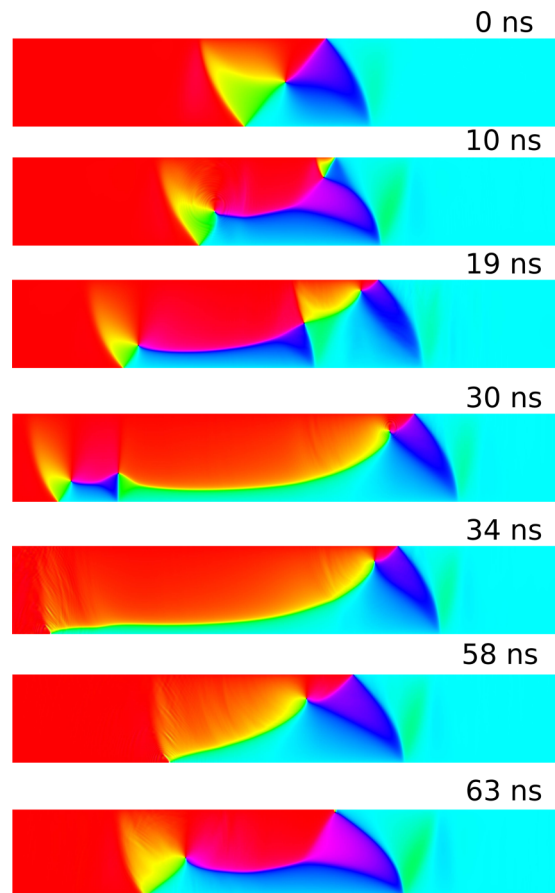


FIG. 9. DW dynamics starting from a VW initial state for $w = 1536$ nm, $\Delta_z = 15$ nm and $B_{\text{ext}} = 2.1$ mT. The DW dynamics shows the nucleation of an antivortex into the strip from the top edge, followed by the injection of a vortex. The antivortex moves towards the vortex that was previously in the strip, leading to an annihilation process between them. After that, the DW structure reduces its width, resulting a VW. Then the process is repeated.

processes are observed, such as one involving a vortex and two antivortices.

For $2.2 \leq B_{\text{ext}} \leq 3.6$ mT, an extended region where the system is again unable to support compact DWs is encountered. Two different types of instability have been observed. For $2.2 \leq B_{\text{ext}} \leq 3.4$ mT the instability is like the one shown in Fig. 6(b), where the DW width grows without limit. For $B_{\text{ext}} = 3.5$ and 3.6 mT, although the DW width also increases without limit, at the same time several vortex-antivortex annihilation processes take place.

For even larger fields, compact DWs are observed again, with their large velocity increasing linearly with B_{ext} . Notice that for $B_{\text{ext}} = 4$ mT, v exceeds 550 m/s. In this regime, the DW dynamics is quite chaotic with transformations between complex dynamical states. In Fig. 10 the DW dynamics is shown for $B_{\text{ext}} = 3.8$ mT (see also Ref. [49], movie 6). Repeated annihilation processes between vortices and antivortices are observed. Note that in this regime, creation of vortices and antivortices occurs not only at the edges (as for lower fields), but also within the strip (see snapshots at 6, 12, and 21 ns in Fig. 10). Similar processes of

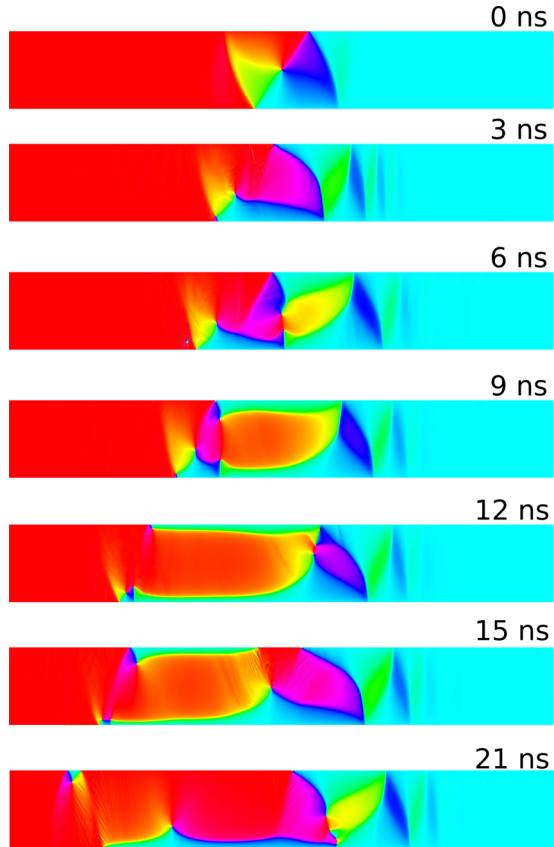


FIG. 10. DW dynamics starting from a VW initial state for $w = 1536$ nm, $\Delta_z = 15$ nm and $B_{\text{ext}} = 3.8$ mT. Very complex, possibly chaotic dynamics is observed.

vortex-antivortex pair generation within the strip have been reported before [56].

We have also analyzed the dependence of the VW velocity within the low-field steady/viscous regime as a function of w and Δ_z , within a range of w and Δ_z such that VW is the equilibrium structure. The main panel of Fig. 11 shows the VW velocity as a function of w for $\Delta_z = 15$ nm and $B_{\text{ext}} = 0.3$ mT, displaying an increasing trend of v with w . Since the width of VWs in wider strips is larger, such a trend is in agreement with the 1d model, predicting a linear dependence of v on the DW width [10]. This can also be rationalized by the fact that a given amount of spin rotation leads to DW displacements proportional to the DW width. Inset of Fig. 11 displays v as a function of Δ_z for $w = 768$ nm and $B_{\text{ext}} = 0.3$ mT, showing that the VW velocity decreases with the strip thickness. This may be attributed to the fact that in the steady state, the DW propagation occurs by a precessional motion around an out-of-plane demagnetizing or stray field, which is created by precession in the external field [9]. As the thickness of the strip increases, this demagnetizing field decreases, leading to a decreasing DW velocity.

IV. DOUBLE VORTEX WALL DYNAMICS

For larger strips with widths around $1\mu\text{m}$ and above, the DVW is the equilibrium DW structure [26]. This structure includes two vortices with opposite sense of rotation, see

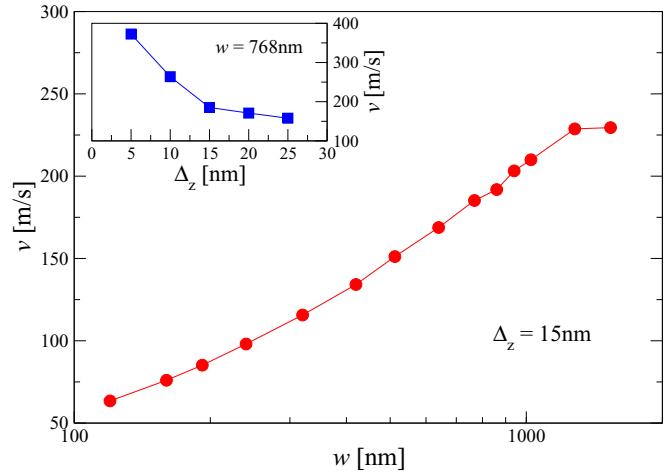


FIG. 11. VW velocity v as a function of the width of the strip w when the thickness and field are fixed to $\Delta_z = 15$ nm and $B_{\text{ext}} = 0.3$ mT, respectively. w is in logarithmic scale. (Inset) VW velocity v as a function of Δ_z for $w = 768$ nm and $B_{\text{ext}} = 0.3$ mT.

Fig. 1(b). As occurs for the VW, the DVW presents different dynamical behaviors depending on the lateral dimensions of the strip, and the magnitude of the driving field. Moreover, in the case of DVW the polarities of the two vortices are an important factor that strongly affects the dynamics. For small enough B_{ext} , the DVW structure is stable depending on the polarities of the vortices. For the case in which the DVW is stable, the DW dynamics exhibits a unique low-field regime. However, the topology of the DVW, consisting of two vortices and four HAV edge defects [54] [see Fig. 12(b)] implies that the energy barrier for transforming the DVW into a VW is not very high. Thus, already for relatively low fields, DVWs tend to transform into VWs, via a process in which one of the vortices annihilates with two edge defects at the edge of the strip. Since the total winding number of a vortex and two edge defects is zero, this process happens easily, i.e., it does not require a large applied field.

This can be illustrated by considering, e.g., a strip with $w = 1536$ nm and $\Delta_z = 15$ nm, corresponding again to the phase boundary between VW and DVW equilibrium structures where they have the same energy [26]. The $v(B_{\text{ext}})$ curve obtained by starting from a DVW initial DW with equal polarity for both vortices (parallel case) is represented in Fig. 8, together with the corresponding data for the VW initial state in the same strip geometry (see the previous section). For low fields ($B_{\text{ext}} \leq 0.6$ mT), the $v(B_{\text{ext}})$ curves are quite different for the two structures, indicating that the DVW structure is stable against small applied fields, and thus exhibits a unique small field steady/viscous regime. In this regime, the moving DVW structure initially exhibits transient oscillatory relative motion of the two vortex cores, with the oscillation amplitude decreasing with time and eventually vanishing when the regime with steady DVW motion is reached [see Fig. 13(a) and Ref. [49], movie 7]. In this low-field regime, the VW moves faster than the DVW. This is due to the increased energy dissipation related to the second vortex core of the DVW.

However, if the polarities of the vortices are opposite (antiparallel case), the dynamics at low fields is very different.

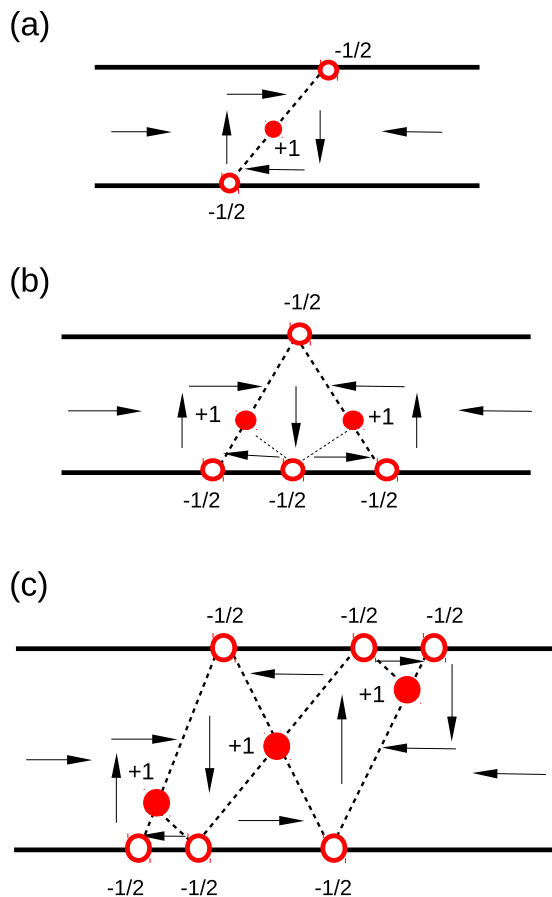


FIG. 12. The equilibrium DW structures in terms of their constituent elementary topological defects, i.e., vortices with a winding number $+1$ (full circles) and edge defects with winding number $-1/2$ (empty circles). (a) Vortex wall, (b) double vortex wall, and (c) triple vortex wall.

DVW is unstable and transforms into a VW, see Fig. 13(b). Contrary to the static case where the polarity barely contributes to the energy of the equilibrium state [26], in the dynamics it plays an important role. This is because when an external field is applied, the vortex core feels a gyrotropic force [48] $\mathbf{F}^g = pG\hat{\mathbf{z}} \times \mathbf{v}$, that depends on the polarity; $G = 2\pi Jt$ is the gyrotropic constant and \mathbf{v} the velocity. In the case of a VW, this force tries to expel the vortex out through the top or bottom edge depending on the polarity, and independently of the vortex chirality. For DVWs with parallel polarity [see the panel top of Fig. 13(a)], both vortices will try to leave the strip through the same edge due to an equal gyrotropic force exerted on them. As in a DW structure the topological defects have to be compensated for, i.e., the total winding number must be equal to zero, only one of the vortices can be expelled. As a result, a competition between them appears, which tends to keep both vortices in the strip, see Fig. 13(a).

In the case of DVW with antiparallel polarity [see top in Fig. 13(b)], the gyrotropic force on one of the vortices is opposite to the one exerted on the other. This means that one vortex will move to the top edge, and the other to the bottom one. As a consequence one of the vortices will be expelled, i.e., the one located closer to the edge in its direction of movement. Figure 13(b) shows how the DVW transforms into a VW.

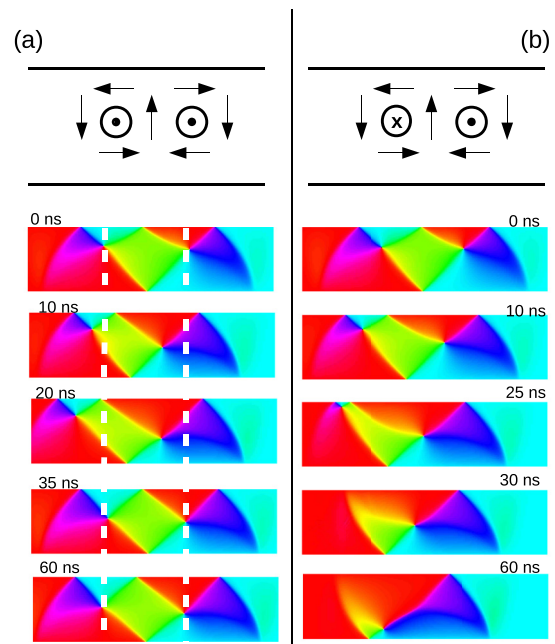


FIG. 13. DVW dynamics in a strip with $w = 1536$ nm and $\Delta_z = 15$ nm for $B_{\text{ext}} = 0.3$ mT and different polarities of the vortices. (a) Parallel polarity; dashed lines help to see the transient relative motion of the vortex cores, before the steady state motion with a nonevolving DVW structure is reached. (b) Antiparallel polarity; DVW is not stable, and becomes a VW.

Returning to the parallel case, for $B_{\text{ext}} = 0.7$ mT, the external field is strong enough to push one of the vortices out of the system (a process that takes place via annihilation of the vortex with the two HAV edge defects), and the DW transforms into a VW. Then, the steady-state velocity equals that of the corresponding VW initial state, a behavior that persists up to $B_{\text{ext}} = 1.1$ mT (see Fig. 8). Thus also here the attraction-repulsion effect is observed. However, the DW dynamics exhibits some dependence on the initial state also for fields exceeding 1 mT. For $B_{\text{ext}} = 1.2$ mT, the DVW initial state leads to a Walker breakdown-type of behavior, with repeated transitions between different DW structures (VW and ATW), similarly to the observations in Fig. 6(a). However, the corresponding system with a VW initial state displays the attraction-repulsion effect, and a significantly larger v for the same field. Also, while for a VW initial state $B_{\text{ext}} = 1.3$ and 1.4 mT lead to stable attraction-repulsion type of behavior, starting from a DVW initial state leads to unstable behavior where the DW width grows without limit. For $B_{\text{ext}} = 1.5$ mT, such an instability is encountered also when starting from the VW initial configuration. For $B_{\text{ext}} = 1.6$ mT, the instability occurs when the VW is the initial state, but not for the case in which DVW is the initial configuration. For larger fields, the two initial conditions lead to the same steady state dynamics, described in the previous section for the VW initial state, see Figs. 9 and 10.

Similar DVW dynamics is observed also for other w and Δ_z values with DVW as the equilibrium structure. However, in the parallel case the “critical” field magnitude at which one of the two vortices is expelled from the strip, leading to a VW, depends on w and Δ_z in a nontrivial fashion. For some

systems belonging to the DVW equilibrium phase, this field is quite small, around 0.1 mT. A possible explanation is that the energy required to move one of the vortices towards the strip edge is smaller in wider strips, given that shape anisotropy (an edge effect) is weaker in wider strips.

V. TRIPLE VORTEX WALL DYNAMICS

For strips with even larger lateral dimensions, TVW is the equilibrium DW structure [26], see Fig. 1(b). As in the case of DVW, also TVW tends to transform into simpler DW structures under applied fields for topological reasons. TVWs are composed of three vortices and six HAV edge defects, see Fig. 12(c). Thus, one or two vortices can be expelled from the strip with relative ease, given that they can annihilate with two HAVs at the strip edge in a process where the total winding number is conserved. Only in the case of very small fields the TVWs maintain their structure during motion depending on the polarities of the vortices. As in the DVW case, the polarity is fundamental for the stability of the TVW structure. As for TVWs there are three different vortices, the situation is more complex than in the case of DVWs. Figure 14 shows the

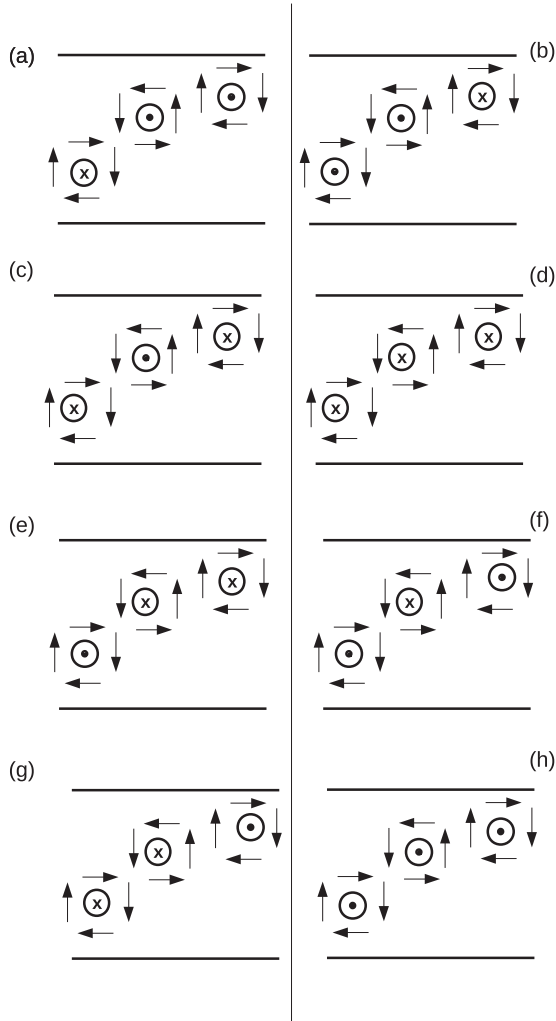


FIG. 14. Top view sketch of the different polarity configurations for a TVW structure.

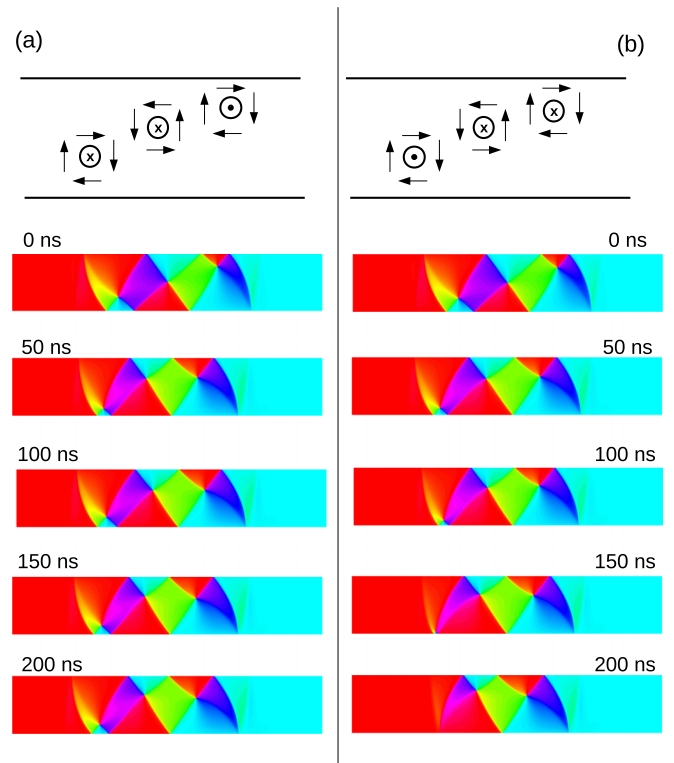


FIG. 15. TVW dynamics in a strip with $w = 5120$ nm and $\Delta_z = 25$ nm for $B_{\text{ext}} = 0.05$ mT and two different configurations of the polarity. (a) TVW structure is stable against this field. (b) The initial TVW structure transforms into a DVW, which then exhibits steady dynamics within the simulation time scale.

different combinations of the polarities of the TVW structure. Contrary to the DVWs, where for parallel polarity of the vortices the DW structure is stable at low fields, for TVW the sign of the polarity is also important. For a strip with $w = 5120$ nm and $\Delta_z = 25$ nm under $B_{\text{ext}} = 0.05$ mT, we have observed that when the three vortices are parallel with polarity $p = -1$ [Fig. 14(d)], the TVW is stable, whereas for polarity $p = +1$ [Fig. 14(h)] TVW transforms into a DVW (not shown). This is related to the gyrotropic force exerted on the vortices and the interactions between them. Depending on the polarity the gyrotropic force acts in one direction or the other, facilitating or preventing the expulsion of the vortex out of the strip.

On the other hand, for antiparallel polarity also the position of the vortices with a given polarity is essential in the TVW dynamics. In Figs. 14(e) and 14(g), two different polarity configurations are shown, where two vortices have polarity $p = -1$ and the third one has $p = +1$. These two configurations show different DW dynamics. To illustrate this, we consider a strip with $w = 5120$ nm and $\Delta_z = 25$ nm, where TVW is one of the equilibrium DW structures [26], considering both polarity configurations. For $B_{\text{ext}} = 0.05$ mT and one of the above mentioned configuration, TVW is stable, keeping its structure during the simulation time, see Fig. 15(a). For the second one, the initial TVW transforms into a DVW as one of the three vortices is expelled from the strip. The resulting DVW exhibits steady dynamics within the simulation time

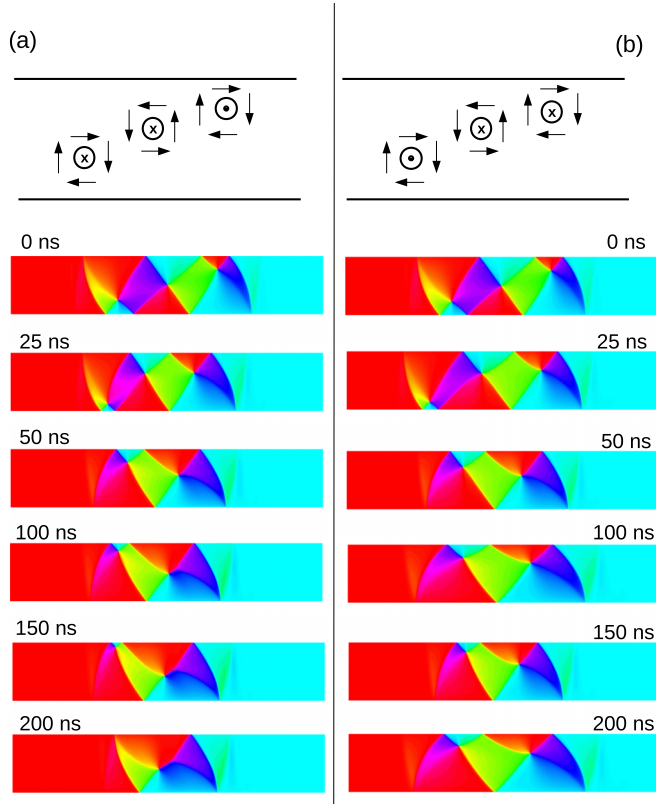


FIG. 16. TVW dynamics in a strip with $w = 5120$ nm and $\Delta_z = 25$ nm for $B_{\text{ext}} = 0.1$ mT and two different configurations of the polarity. (a) The initial TVW structure transforms first into a DVW, and then into a structure exhibiting the topology of a VW. (b) The initial TVW structure transforms into a DVW, which then exhibits steady dynamics within the simulation time scale.

scale, see Fig. 15(b). This is because the vortex expelled from the strip is the one with opposite polarity, and thus the two remaining vortices in the strip are parallel.

To understand why one polarity configuration is stable and the other not, we have to consider the directions of movement of the vortices when an external magnetic field is applied. For the first configuration [Fig. 15(a)], due to the gyrotropic force, the vortices situated at the left and in the center move to the top strip edge, while the third one moves towards the bottom edge. The vortices of the ends of the DW structure are very far from the edges towards which they are moving, and the vortex at the middle cannot be expelled due to the chirality of the vortices. If the vortex at the center was expelled, the two remaining vortices had equal sense of rotation, and the resulting DW structure would not respect the basic principles of topological defects in DWs. To obtain a total winding number equal to zero, new topological defects should be created with the corresponding cost in energy [26]. As a result, the TVW is stable at this magnetic field. For the other case [Fig. 15(b)], the vortex on the left moves to the bottom edge, whereas the other two move towards the top edge. Here, the vortices at the ends are very close to the edges, and in principle we could expect that any of them could be expelled. However, as the vortex at the center also tries to go out, a competition between this vortex and the one situated at the right appears, keeping both inside the strip. Thus the vortex on the left is expelled more easily.

As the applied field increases, the TVW transforms into simpler DW structures for all combinations of the vortex polarities. This can be clearly seen in Fig. 16, where the DW dynamics for the same strip with both polarity configurations are represented for $B_{\text{ext}} = 0.1$ mT. For the configuration where the TVW is stable at a smaller field, now two of the three vortices are expelled, leading to a VW, see Fig. 16(a). Here the applied field is strong enough to push one of the vortices (the left one) out of the strip. The two remaining vortices have opposite polarities. As it was already seen in the previous section, the DVW with antiparallel polarity is unstable, and the DW structure transforms into a VW. Figure 16(b) shows that for the other configuration, the DW dynamics is similar to the one observed for $B_{\text{ext}} = 0.05$ mT. TVW transforms in a

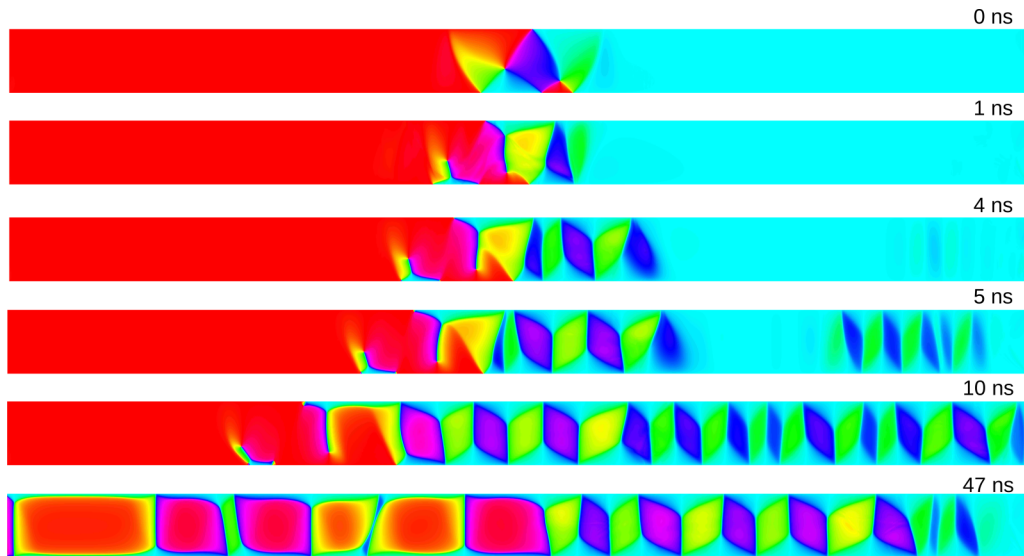


FIG. 17. Snapshots from a strip with $w = 5120$ nm and $\Delta_z = 5$ nm with $B_{\text{ext}} = 0.9$ mT. For these conditions the system cannot support a compact DW, resulting in an instability where a complex magnetization pattern fills the strip.

DVW, which is stable during the simulation time because the polarities of the two vortices are parallel.

Similar results have been found for a larger strip with $w = 6144$ nm and $\Delta_z = 25$ nm (not shown). For this case, once two of the three vortices have been expelled, the remaining DW has the topology of a VW. However, this structure is more complex than typical VWs observed in more narrow strips.

Finally, we mention the fact that also for these very wide strips the system is, for some combinations of the strip geometry and the applied field, unable to support a compact DW. The related instabilities can be quite complex, as illustrated in Fig. 17 for a strip with $w = 5120$ nm, $\Delta_z = 5$ (corresponding to the phase boundary between VW and DVW equilibrium structures [26]), and $B_{\text{ext}} = 0.9$ mT. In this case, a complex domain/magnetization pattern starts to emanate from the initial DW, eventually filling the entire strip. Thus, for such wide permalloy strips, one starts to reach a situation resembling increasingly that in thin films rather than nanostrips, where system-spanning closure domain/Landau patterns are typically observed [57–59].

VI. SUMMARY AND CONCLUSIONS

To summarize, we have analyzed in detail DW dynamics driven by an external magnetic field in wide permalloy strips, by considering a set of example geometries spanning a wide range of strip widths. Our main message is that the system geometry crucially affects the DW dynamics also for lateral strip dimensions beyond the typical nanostrip ones which have been extensively studied in the literature. Already for relatively confined systems with VW as the stable DW structure, several different dynamical behaviors are obtained, including processes such as various periodic transitions between different DW structures, the attraction-repulsion effect, etc. These behaviors are also visible in the profiles of the resulting $v(B_{\text{ext}})$ curves. For wider strips where the DVWs and TVWs are the equilibrium DWs, the polarities of the vortices play an important role in the dynamics. For small fields, DVW and TVW structures are stable, i.e., they maintain their structure depending on the polarity and the interactions between the vortices. As the size of the strip increases, for relatively small fields in the range of a few millitesla, these structures tend to transform towards DW structures with less

vortices. This effect is intimately linked to the topology of such more complex equilibrium DW structures. The latter can be represented as composite objects of topological defects, consisting of vortices and half-antivortex edge defects [54]. The transitions towards simpler DW structures proceed via annihilation processes conserving the winding number, and thus the energy barrier for such transitions tends to be low.

Another key feature of wide permalloy strips is that they are not always able to support compact DWs. Indeed, as the strip width increases, shape anisotropy becomes weaker, eventually leading to a situation where increasing the DW width without limit becomes energetically favorable. In dynamical situations addressed here this tends to happen via different DW components acquiring different field-driven propagation velocities [14], leading to gaps in the $v(B_{\text{ext}})$ curves as the DW and therefore its velocity become ill-defined concepts. More generally, as the strip width increases, one would expect a crossover to a thin filmlike behavior with system-spanning closure domain/Landau patterns [57–59] instead of localized DWs.

From an experimental perspective, our study highlights the importance of performing simultaneous analysis, e.g., by micromagnetic simulations, to verify which of the large variety of complex dynamical behaviors corresponds to the experimental conditions at hand. While our preliminary studies suggest that the attraction-repulsion effect is stable against adding random structural disorder [60] of moderate strength, future studies should address the effects of, e.g., thermal fluctuations and quenched disorder on DW dynamics in wide permalloy strips. For instance, the effect of e.g. edge roughness on the periodic vortex core dynamics near the strip edge would be interesting to address in detail.

ACKNOWLEDGMENTS

We thank Ilari Rissanen for sharing computational resources. This work has been supported by the Academy of Finland through its Centres of Excellence Programme (2012–2017) under Project No. 251748, and via an Academy Research Fellowship (LL, Project No. 268302). We acknowledge the computational resources provided by the Aalto University School of Science “Science-IT” project, as well as those provided by CSC (Finland).

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